



# An empirical spectroscopic model for eleven electronic states of VO

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## ABSTRACT

Previously-determined empirical energy levels are used to construct a rovibronic model for the  $X^4\Sigma^-$ ,  $A'^4\Phi$ ,  $A^4\Pi$ ,  $B^4\Pi$ ,  $C^4\Sigma^-$ ,  $D^4\Delta$ ,  $1^2\Delta$ ,  $1^2\Sigma^+$ ,  $1^2\Phi$ ,  $1^2\Pi$  and  $2^2\Pi$  electronic states of vanadium monoxide. The spectrum of VO is characterized by many couplings and crossings between the states associated with these curves. The model is based on the use of potential energy curves, represented as extended Morse oscillators, and couplings (spin-orbit, spin-spin, angular momentum), represented by polynomials, which are tuned to the data plus an empirical allowance for spin-rotation couplings. The curves are used as input for the variational nuclear motion code Duo. For the  $X^4\Sigma^-$ ,  $1^2\Phi$  and  $1^2\Pi$  states the model reproduces the observed energy to significantly better than  $0.1\text{ cm}^{-1}$ . For the other states the standard deviations are between  $0.25$  and  $1.5\text{ cm}^{-1}$ . Further experimental data and consideration of hyperfine effects would probably be needed to significantly improve this situation.

## 1. Introduction

Vanadium monoxide (VO) is a diatomic molecule with a complicated open shell structure [1,2]; it absorbs strongly at red and near infrared wavelengths. VO is thought to provide an important source of opacity in the atmospheres of both cool stars [3] and of exoplanets [4,5]. While VO bands are well-known in the spectrum of cool stars [6] and sunspots [7], only tentative detections have been made in the atmospheres of exoplanets [8–14]. The charged species  $\text{VO}^-$  and  $\text{VO}^+$  are also important targets of theoretical and experimental studies, see e.g., Miliordos and Mavridis [15] and the references therein.

Theoretically modeling of VO spectra is a challenge both from the perspective of electronic structure calculations [15–18] and, as will be seen below, from the perspective of representing the nuclear motion due to many resonances between different states. McKemmish et al. [19] constructed the so-called VOMYT line list for VO as part of the ExoMol project [20,21]. The VOMYT line list probably provides a reasonable basis for modeling VO opacities but has been demonstrated to be inadequate for high resolution studies of exoplanetary spectra [22].

The spectrum of VO has been well-studied experimentally [23–35], see the summary by Bowersman et al. [36]; however, the complexity of its spectrum means that there are still important gaps in these studies. Although, as discussed below, our general approach to constructing a spectroscopic model for open shell molecules such as VO is based on variational solutions of the nuclear motion Schrödinger equation, our work on VO [2,36] and other open shell molecules [37–40] has

made extensive use of Western's effective Hamiltonian program PGO-PHER [41]. We favor the method based on variational solutions as it usually shows much more stable extrapolation properties which is an important feature when one wishes to compute line lists for hot molecules.

In this paper, we attempt to build a spectroscopic model which covers the most important low-lying states of VO, including 6 quartet and 5 doublet electronic states; given the difficulty in fully characterizing the couplings between these state we limit ourselves to considering only the coupling terms necessary to provide a reasonable spectroscopic model. The 11 electronic states considered are shown in Fig. 1; these state are all part of the VOMYT model. They are responsible for the main absorption bands in the near infrared and visible, or are states which are known to interact with them.

The purpose of this work is to provide the basis for a more reliable model than VOMYT as well providing for the quantum number assignments for unassigned recent observations [35,36]. Thus, our focus is on the accuracy of energy levels and we develop an empirical model that excludes most of the off-diagonal couplings considered in VOMYT [19]. In the recent study by Bowersman et al. [36], a spectroscopic network of VO transitions was constructed as a prelude to a MARVEL (measured active rotation vibration energy levels) [42] study of the 11 states involved in this network. The next two sections will demonstrate our effort to reproduce the rovibronic energy levels obtained from the MARVEL study. The other two states ( $a^2\Sigma^-$  and  $b^2\Pi$ ) included in VOMYT [19] were not analyzed in the work of Bowersman et al. [36] and are, therefore, excluded from our model as well.

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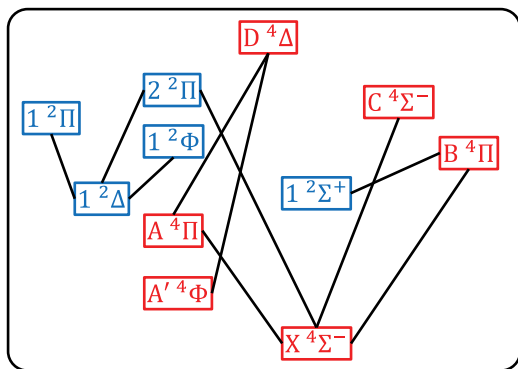


Fig. 1. Band systems connecting the 11 electronic states included in the model of this work. See Bowersman et al. [36] for a more comprehensive diagram.

Table 1

Vibrational states considered for each electronic state. Table 5 of Bowersman et al. [36] lists the corresponding experimental studies.

Number	States	$v$
1	$X^4\Sigma^-$	0-2
2	$A'^4\Phi$	0-2
3	$A^4\Pi$	0
4	$B^4\Pi$	0-1
5	$C^4\Sigma^-$	0-2
6	$D^4\Delta$	0-1
7	$1^2\Delta$	0-1
8	$1^2\Sigma^+$	2-3
9	$1^2\Phi$	0
10	$1^2\Pi$	0-3
11	$2^2\Pi$	0-1

## 2. Spectroscopic model

The model was developed using the open-source program for variational calculation of diatomic spectra, Duo<sup>1</sup> [43]. All nuclear motion calculations used 10 vibrationally contracted basis functions for each electronic state based on 401 sinc-DVR grid points, covering the internuclear distance range from 1.2 to 4 Å. The upper bound of the energy was set to 50 000 cm<sup>-1</sup>, however energy levels of interest for this work are expected to be below 21 000 cm<sup>-1</sup>. The number of vibrational levels considered for each electronic state is listed in Table 1.

### 2.1. Potential energy curves

The potential energy curves (PECs) of all states are represented by the second-order extended Morse oscillator (EMO) function [44]

$$V(R) = T_e + (A_e - T_e) \left[ 1 - \exp(-\beta_{\text{EMO}}(R)(R - R_e)) \right]^2, \quad (1)$$

where  $R$  is the internuclear distance,  $R_e$  is the equilibrium bond length,  $T_e$  is the potential minimum and  $A_e$  is the corresponding asymptote;  $\beta_{\text{EMO}}$  is given by the formula:

$$\beta_{\text{EMO}}(R) = b_0 + b_1 y_p(R) + b_2 y_p^2(R), \quad (2)$$

where  $y_p(R)$  is:

$$y_p(R) = \frac{R^p - (R_e)^p}{R^p + (R_e)^p}. \quad (3)$$

Duo's users can set independent integer values of  $p$  on the left- and right-hand sides of  $R_e$  and we chose  $p = 4$  for both sides. The parameters of fitted EMO functions are listed in Table 2.

### 2.2. Diagonal couplings within a electronic state

All coupling curves are represented by low-order polynomials

$$p(R) = a_0 + \sum_{n=1}^{\infty} a_n (R - R_e)^n. \quad (4)$$

Our aim is to use the fewest terms to satisfactorily reproduce the experimental energy levels. Thus, we started with uncoupled PECs for each state and then, added those coupling curves necessary to reduce the fitting errors.

We assume the spin splittings of  $X^4\Sigma^-$  and  $C^4\Sigma^-$  mainly depend on spin–spin coupling terms. Although the assumption is not supported by *ab initio* calculations, e.g. see our previous discussion of the  $X^4\Sigma^-$  state [2], this empirical representation can usually give accurate energy levels. Apart from the spin–spin interaction, spin–rotation terms were used to resolve  $J$ -dependent spin splittings of  $^2\Sigma$  and  $^4\Sigma$  states, where  $J$  is the total angular momentum.

In contrast, spin–orbit terms dominate the spin splittings of  $\Pi$ ,  $\Delta$ ,  $\Phi$  and  $\Gamma$  states and are  $J$ -dependent too. Therefore, we replaced spin–rotation curves with spin–orbit coupling ones for those states. Spin–spin coupling curves were still used for the quartet states, i.e.  $A'^4\Phi$ ,  $A^4\Pi$ ,  $B^4\Pi$  and  $D^4\Delta$ .

The optimized parameters of diagonal spin–spin, spin–rotation and spin–orbit couplings are listed in Tables 3–5, respectively.

We also modeled the  $A$ -doubling of  $\Pi$  states by using the empirical terms defined by Brown and Merer [45]:

$$\langle \mp 1, \Sigma \pm 2, J, \Omega | \hat{H}_{\text{LD}} | \pm 1, \Sigma, J, \Omega \rangle = \frac{1}{2} (o_v + p_v + q_v) \times \sqrt{[S(S+1) - \Sigma(\Sigma \pm 1)][S(S+1) - (\Sigma \pm 1)(\Sigma \pm 2)]}, \quad (5)$$

$$\langle \mp 1, \Sigma \pm 1, J, \Omega \mp 1 | \hat{H}_{\text{LD}} | \pm 1, \Sigma, J, \Omega \rangle = -\frac{1}{2} (p_v + 2q_v) \times \sqrt{[S(S+1) - \Sigma(\Sigma \pm 1)][J(J+1) - \Omega(\Omega \mp 1)]}, \quad (6)$$

$$\langle \mp 1, \Sigma, J, \Omega \mp 2 | \hat{H}_{\text{LD}} | \pm 1, \Sigma, J, \Omega \rangle = \frac{1}{2} q_v \times \sqrt{[J(J+1) - \Omega(\Omega \mp 1)][J(J+1) - (\Omega \mp 1)(\Omega \mp 2)]}. \quad (7)$$

In these equations,  $o$ ,  $p$  and  $q$  depend on vibrational quantum number,  $v$ , as Brown and Merer derived the equations based on effective Hamiltonians. In our variational model, we use coupling curves instead. For quartet states, we include all the three coupling curves, while for doublet states, Eq. (5) is automatically zero and we only need the curves of  $p+2q$  and  $q$ . The fitted parameters for the  $o+p+q$ ,  $p+2q$  and  $q$  curves are listed in Table 6.

#### 2.2.1. Off-diagonal spin–orbit couplings

We can observe obvious avoided-crossing structures from the rovibronic energy distributions of  $B^4\Pi - 1^2\Sigma^+$  and  $C^4\Sigma^- - 2^2\Pi$  coupled states. Thus, off-diagonal spin–orbit coupling curves, whose parameters are given in Table 7, were introduced to model those structures in crossing regions. The  $B^4\Pi - 1^2\Sigma^+$  spin–orbit coupling term was optimized using Duo, while the  $C^4\Sigma^- - 2^2\Pi$  coupling was manually given an estimated value. See the next section where the results are discussed.

## 3. Results

### 3.1. Uncoupled states

The fitting residues of the uncoupled quartet states are given in Fig. 2. The energy levels of  $X^4\Sigma^-$  and  $A^4\Pi$  are accurately reproduced. Compared with the  $X^4\Sigma^-$  and  $A^4\Pi$  states, the spin-splittings of  $A'^4\Phi$  and  $D^4\Delta$  are not well estimated which may arise from electronic state interactions not included in our model. However, the assignment will not be affected as the energy gaps due to spin–orbit couplings of  $A'^4\Phi$  and  $D^4\Delta$  are usually more than a hundred wavenumbers.

The fitting residues of the uncoupled doublet states are shown in Fig. 3. The doublet states only have two fine structure series which

<sup>1</sup> <https://github.com/exomol/Duo>

**Table 2**  
Optimized extended Morse oscillator (EMO) parameters for the potential energy curves, see Eq.(1).

Number	State	$T_e$ [cm <sup>-1</sup> ]	$R_e$ [Å]	$b_0$ [Å <sup>-1</sup> ]	$b_1$ [Å <sup>-1</sup> ]	$b_2$ [Å <sup>-1</sup> ]
1	X <sup>4</sup> Σ <sup>-</sup>	0	1.58947260	1.87183393	6.042 606 63 × 10 <sup>-3</sup>	-1.158 731 02 × 10 <sup>-1</sup>
2	A' <sup>4</sup> Φ	7292.4851	1.62591915	1.88461684	-1.332 091 48 × 10 <sup>-2</sup>	1.585 203 16 × 10 <sup>-1</sup>
3	A <sup>4</sup> Π	9534.8342	1.63405670	1.92275217		
4	B <sup>4</sup> Π	12659.4107	1.64032109	1.93810207	0	-2.273 213 32 × 10 <sup>-1</sup>
5	C <sup>4</sup> Σ <sup>-</sup>	17494.4977	1.67098029	1.94981612	-5.170 578 88 × 10 <sup>-2</sup>	3.271 206 99 × 10 <sup>-1</sup>
6	D <sup>4</sup> Δ	19238.8720	1.68337218	1.91485252	0	3.03001473
7	1 <sup>2</sup> Δ	9867.6920	1.58225344	2.06276090	0	-4.155 158 02 × 10 <sup>-1</sup>
8	1 <sup>2</sup> Σ <sup>+</sup>	10376.3396	1.57856000	2.15102878	0	-3.663 042 53 × 10 <sup>-1</sup>
9	1 <sup>2</sup> Φ	15442.9452	1.62904530	2.09925561		
10	1 <sup>2</sup> Π	17099.9353	1.62955854	2.10656122	2.732 894 47 × 10 <sup>-1</sup>	1.641 222 98 × 10 <sup>-1</sup>
11	2 <sup>2</sup> Π	18139.4421	1.62244306	2.15748234		

**Table 3**  
Optimized polynomial parameters for diagonal spin-spin coupling curves.

State	$a_0$ [cm <sup>-1</sup> ]	$a_1$ [cm <sup>-1</sup> Å <sup>-1</sup> ]	$a_2$ [cm <sup>-1</sup> Å <sup>-2</sup> ]
X <sup>4</sup> Σ <sup>-</sup>	2.05042406	2.76918955	
A' <sup>4</sup> Φ	-1.08746355		
A <sup>4</sup> Π	1.89104909		
B <sup>4</sup> Π	2.62698261		
C <sup>4</sup> Σ <sup>-</sup>	0.73293575		
D <sup>4</sup> Δ	1.29783687	0	-6.220 435 54 × 10 <sup>2</sup>

**Table 4**  
Optimized polynomial parameters for diagonal spin-rotation coupling curves.

State	$a_0$ [cm <sup>-1</sup> ]	$a_1$ [cm <sup>-1</sup> Å <sup>-1</sup> ]
X <sup>4</sup> Σ <sup>-</sup>	2.143 878 41 × 10 <sup>-2</sup>	1.587 127 88 × 10 <sup>-1</sup>
C <sup>4</sup> Σ <sup>-</sup>	-1.834 643 08 × 10 <sup>-2</sup>	
1 <sup>2</sup> Σ <sup>+</sup>	-7.900 000 00 × 10 <sup>-3</sup>	

makes it easier to get good refinement results. The errors for the high- $J$  levels of 1<sup>2</sup>Δ increase dramatically which is also due to state interactions.

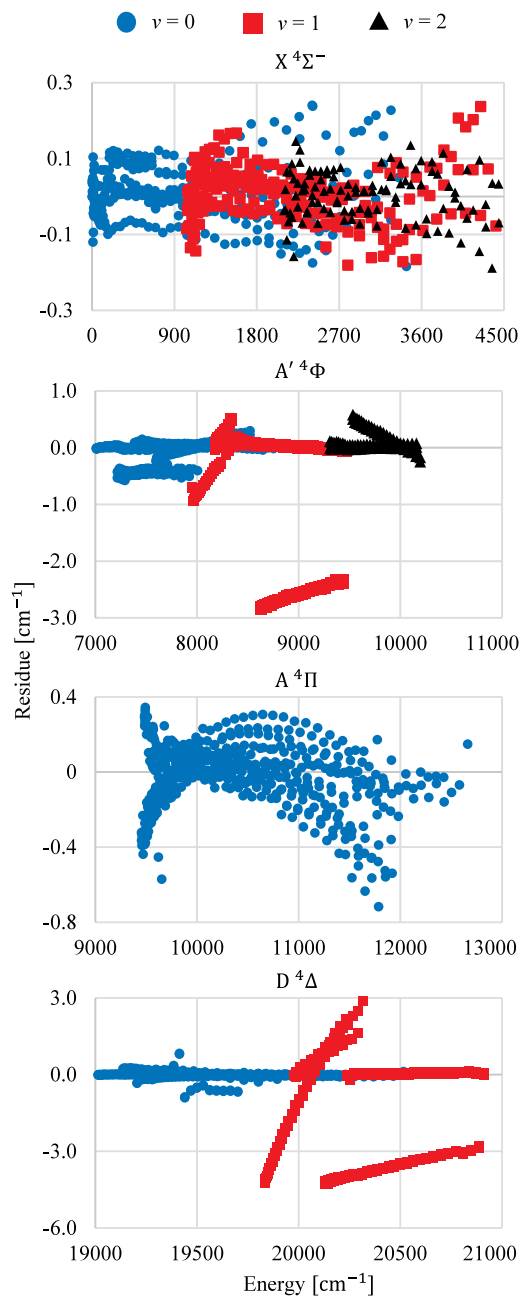
### 3.2. Coupled states

In our model, the systems B<sup>4</sup>Π – 1<sup>2</sup>Σ<sup>+</sup> and C<sup>4</sup>Σ<sup>-</sup> – 2<sup>2</sup>Π are coupled by the corresponding spin-orbit terms to give reasonable avoided crossing structures, as shown in Fig. 4 to Fig. 7.

Figs. 4 and 5 demonstrate the smooth transitions from an electronic state to another in both  $e$  and  $f$  series. Note that, the electronic state quantum numbers of the energy levels near the interaction region are not well-determined. They are assigned according to the dominant contribution of the corresponding basis function to the rovibronic wavefunction. The MARVEL energy levels indicated by the black arrows in these figures were excluded from the fits. They seem to belong to another electronic state but were not recognized in the current model.

The off-diagonal C<sup>4</sup>Σ<sup>-</sup> – 2<sup>2</sup>Π spin-orbit coupling term was not fitted but visually estimated to give an acceptable avoided-crossing structure as shown in Figs. 6 and 7. The  $v = 0, F_2$  levels of 2<sup>2</sup>Π shown in Fig. 6 are recently assigned in the work of Bowesman et al. [36]. However, the crossing point of this series with  $F_2$  series of C<sup>4</sup>Σ<sup>-</sup> is different from the previously determined one, as indicated by the black arrow in Fig. 6. All transitions connecting the C<sup>4</sup>Σ<sup>-</sup> and 2<sup>2</sup>Π states in the MARVEL spectroscopic network are under review and some of them may be reassigned. Thus, we did not make too much effort to improve the accuracy here.

The  $v = 0$  levels of C<sup>4</sup>Σ<sup>-</sup> do not interact with the 2<sup>2</sup>Π state. Thus, the diagonal spin-rotation coupling term in Table 4 was determined by the states with the  $v = 0$  levels. Nevertheless, there are still systematic errors as shown in Fig. 8. The  $J$ -dependent fitting residues of high- $J$  levels should be attributed to C<sup>4</sup>Σ<sup>-</sup> – 1<sup>2</sup>Π interaction, although more experimental data are needed to verify this.



**Fig. 2.** Fitting residues of four uncoupled quartet states: X<sup>4</sup>Σ<sup>-</sup>, A'<sup>4</sup>Φ, A<sup>4</sup>Π and D<sup>4</sup>Δ.

**Table 5**  
Optimized polynomial parameters for diagonal spin-orbit coupling curves.

State	$a_0$ [ $\text{cm}^{-1}$ ]	$a_1$ [ $\text{cm}^{-1} \text{\AA}^{-1}$ ]	$a_2$ [ $\text{cm}^{-1} \text{\AA}^{-2}$ ]	$a_3$ [ $\text{cm}^{-1} \text{\AA}^{-3}$ ]
$A'^4\Phi$	$2.60752743 \times 10^2$	$-3.08377540 \times 10^2$	$-3.01987104 \times 10^3$	$5.20656186 \times 10^4$
$A^4\Pi$	$5.26495455 \times 10^1$			
$B^4\Pi$	$9.70343261 \times 10^1$	$-3.31204551 \times 10^1$		
$D^4\Delta$	$1.39054588 \times 10^2$	0	$2.59920512 \times 10^3$	
$1^2\Delta$	$1.80055832 \times 10^2$	$-1.58772300 \times 10^2$		
$1^2\Phi$	$1.86052851 \times 10^2$			
$1^2\Pi$	$1.24587433 \times 10^2$	$8.26877879 \times 10^1$	$-2.70479601 \times 10^2$	
$2^2\Pi$	$6.03714458 \times 10^1$			

**Table 6**  
Optimized polynomial parameters for the  $o+p+q$ ,  $p+2q$  and  $q$  curves due to  $A$ -doubling coupling.

Coupling	State	$a_0$ [ $\text{cm}^{-1}$ ]
$o+p+q$	$A^4\Pi$	2.05253781
	$B^4\Pi$	1.20113529
$p+2q$	$A^4\Pi$	$-1.41692302 \times 10^{-2}$
	$B^4\Pi$	$3.18974836 \times 10^{-2}$
	$1^2\Pi$	$-3.12107139 \times 10^{-2}$
	$2^2\Pi$	$-2.19602031 \times 10^{-2}$
$q$	$A^4\Pi$	$-2.09279299 \times 10^{-4}$
	$B^4\Pi$	$-1.63337099 \times 10^{-4}$
	$1^2\Pi$	$-2.04345034 \times 10^{-4}$
	$2^2\Pi$	$1.00240289 \times 10^{-3}$

**Table 7**  
Polynomial parameters for the off-diagonal spin-orbit coupling curves.

State	$a_0$ [ $\text{cm}^{-1}$ ]
$B^4\Pi - 1^2\Sigma^+$	21.5107379
$C^4\Sigma^- - 2^2\Pi$	10

**Table 8**  
RMSE and maximum absolute error of each state in  $\text{cm}^{-1}$ .

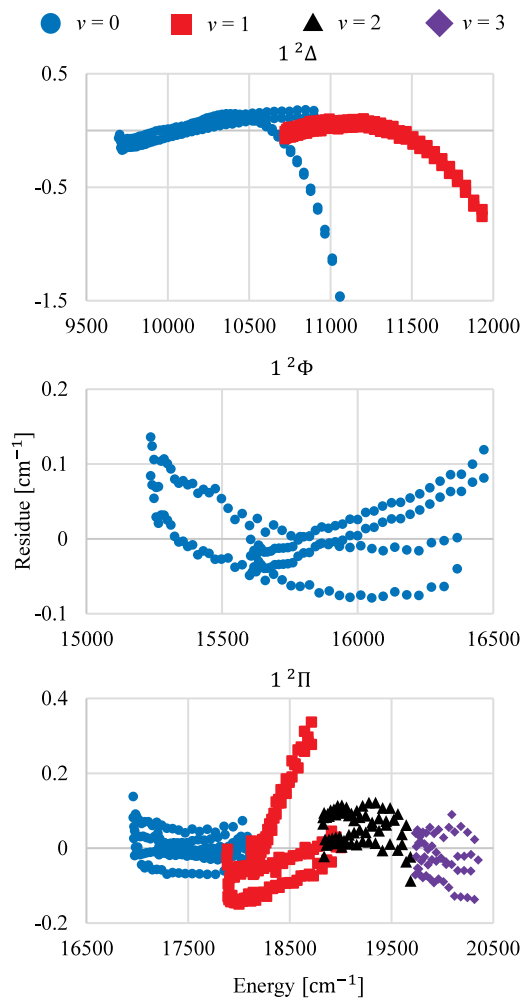
State	$v$	RMSE	Max absolute error
$X^4\Sigma^-$	0	0.0777	0.2404
	1	0.0168	0.2369
	2	0.0651	0.1884
$A'^4\Phi$	0	0.2241	0.5760
	1	1.3634	2.8620
	2	0.206	0.5821
$A^4\Pi$	0	0.1766	0.5131
$D^4\Delta$	0	0.1893	0.9044
	1	2.3556	4.2883
$1^2\Delta$	0	0.2675	1.4664
	1	0.2107	0.7573
$1^2\Phi$	0	0.0507	0.1362
	0	0.0363	0.1380
	1	0.1145	0.3368
	2	0.0696	0.1217
$1^2\Pi$	3	0.0603	0.1372
	0	0.4234	1.2228
	0	0.2949	1.9393
$C^4\Sigma^- - 2^2\Pi$		1.4829	9.0335

### 3.3. Fitting quality

The root-mean-square error (RMSE) of the fitting residues for each electronic state is calculated as

$$\text{RMSE} = \sqrt{\frac{1}{N} \sum_{i=1}^N e_i^2}, \quad (8)$$

where  $e_i$  is the fitting error of the  $i$ th energy level and  $N$  is the number of rovibronic states included. The values of RMSE are listed in Table 8.



**Fig. 3.** Fitting residues of three uncoupled doublet states:  $1^2\Delta$ ,  $1^2\Phi$  and  $1^2\Pi$ .

The VOMYT model was refined using Model Hamiltonian, combination difference energy levels and transition frequencies of VO [19]. As most of the transitions connect to the ground electronic state, VOMYT has a reasonable representation of  $X^4\Sigma^-$ . However, its vibrational structure was still not accurately determined as shown in Fig. 9. There is an obvious offset of the  $v = 2$  energy levels, which is likely to be significantly worse for higher vibrationally excited states. The RMSE value for the VOMYT model is  $0.3015 \text{ cm}^{-1}$ , about four times worse than ours.

The rotational constants of all states in our model are around  $0.5 \text{ cm}^{-1}$  and the rotational energy gaps are estimated as

$$\Delta E = 0.5(J+2)(J+1) - 0.5(J+1)J = J+1 \text{ cm}^{-1}. \quad (9)$$

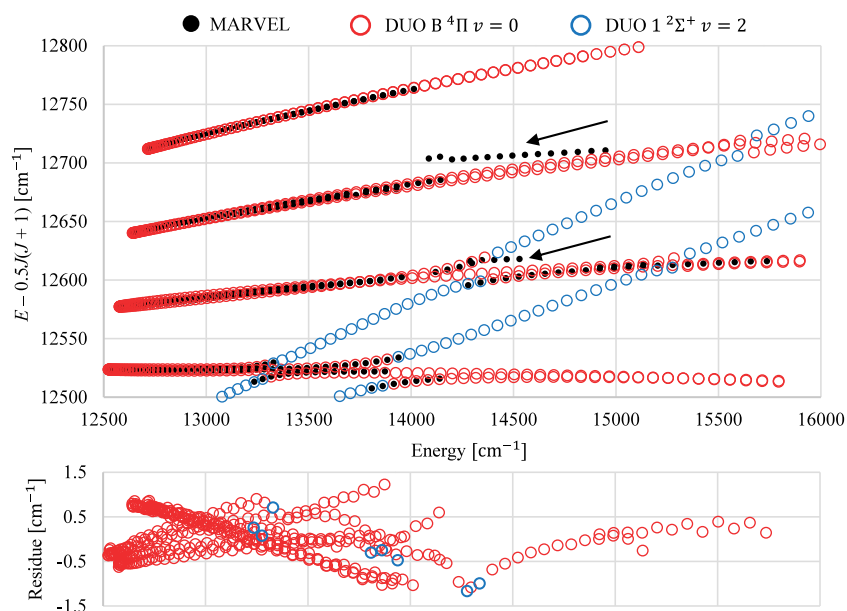


Fig. 4. State interaction of  $B^4\Pi, \nu=0$  and  $1^2\Sigma^+, \nu=2$ . The black dots in the top panel are the rovibronic energy levels obtained from MARVEL analysis [36] while the circles are the fitting results of DUO. The black arrows indicate some unreasonable data which were excluded in our fitting. The bottom panel shows the fitting residues.

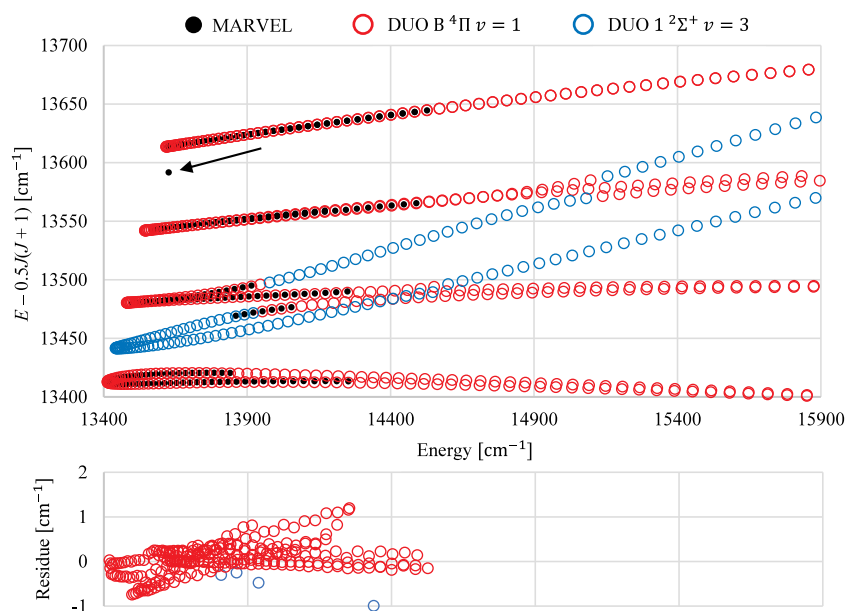


Fig. 5. State interaction of  $B^4\Pi, \nu=1$  and  $1^2\Sigma^+, \nu=3$ . The black dots in the top panel are the rovibronic energy levels obtained from MARVEL analysis [36] while the circles are the fitting results of DUO. The bottom panel shows the fitting residues.

Thus, our model achieves sufficient accuracy to make rotational assignments. The coupling curves can be further refined once revised or extended data is available.

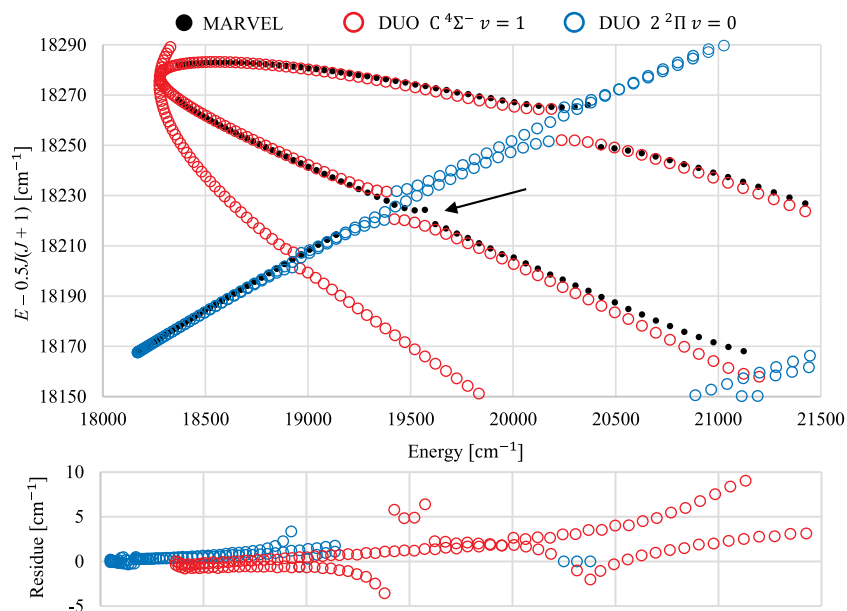
#### 4. Conclusions

We propose an empirical spectroscopic model for 6 quartet and 5 doublet states of VO. The most recent MARVEL analysis of VO [36] generates accurate energy term values for this molecule, which makes it possible to achieve experimental accuracy for some vibronic levels compared with VOMYT [19]. The interactions of  $B^4\Pi - 1^2\Sigma^+$  and

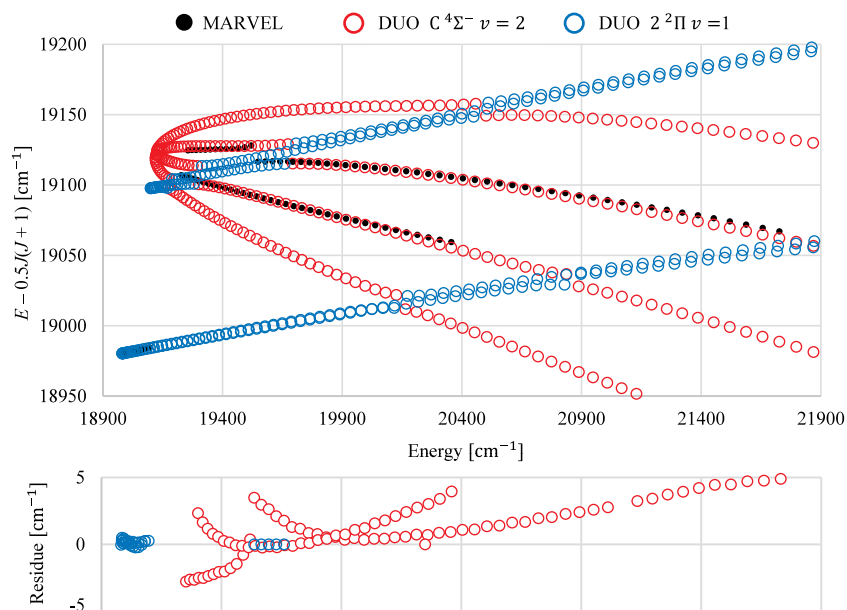
$C^4\Sigma^- - 1^2\Pi$  are analyzed by introducing the off-diagonal spin-orbit couplings between them.

While our model represents a significant advance on what is currently available, there remains two outstanding issues. The first is the treatment of resonance interactions which is discussed extensively above. Given that state-of-the-art *ab initio* calculations on systems with a transition metal atom are not capable of even approaching the accuracy required for high resolution studies, improving our model will require further experimental input.

The second issue is the hyperfine structure. The main isotope of vanadium,  $^{51}\text{V}$ , has a nuclear spin  $I = 7/2$  which is known to lead to pronounced hyperfine effects in the spectrum of VO [25,30,34]. We



**Fig. 6.** State interaction of  $C^4\Sigma^-, v=1$  and  $2^2\Pi, v=0$ . The black dots in the top panel are the rovibronic energy levels obtained from MARVEL analysis [36] while the circles are the fitting results of DUO. The rotational constant of  $2^2\Pi, v=0$  in derived from our model is consistent with that of the recently assigned results (see the blue circles in the top panel). However, its crossing point with the  $F_2$  series of  $C^4\Sigma^-, v=1$  is inconsistent with experimental values (see the area pointed by the black arrow). The bottom panel shows the fitting residues.



**Fig. 7.** State interaction of  $C^4\Sigma^-, v=2$  and  $2^2\Pi, v=1$ . The black dots in the top panel are the rovibronic energy levels obtained from MARVEL analysis [36] while the circles are the fitting results of DUO. The bottom panel shows the fitting residues.

have started to address this issue by extending our nuclear motion program DUO to explicitly include hyperfine effects [46] within its variational model. We have also built a hyperfine-resolved spectroscopic model for the ground  $X^4\Sigma^-$  state of  $^{51}\text{V}^{16}\text{O}$  [2]. We are working on extending the model to other electronic states of VO to give a spectroscopically-accurate, hyperfine-resolved model of the molecule. Given the limited experimental hyperfine-resolved data available for VO, particularly for spectra involving excited vibrational states, further experimental input into this problem would be particularly welcome.

#### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### Data availability

Data is supplied as a supplementary file to the article

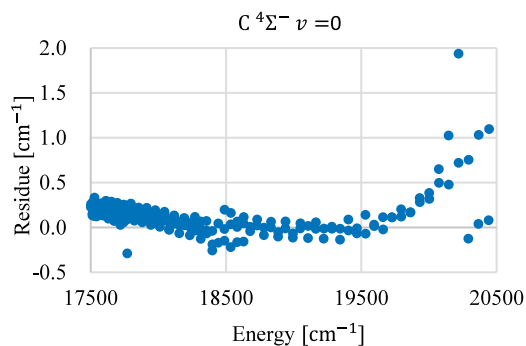


Fig. 8. Fitting residues of  $C^4\Sigma^-, v=0$ .

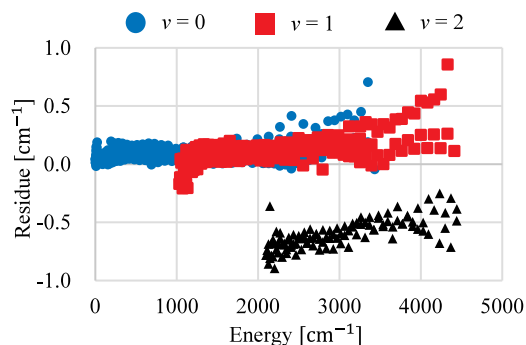


Fig. 9. Energy errors for the  $X^4\Sigma^-$  state as given by the VOMYT model.

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## Appendix A. Supplementary data

Supplementary material related to this article can be found online at <https://doi.org/10.1016/j.jms.2022.111733>.

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