

ExoMol line lists – XLVII. Rovibronic molecular line list of the calcium monohydroxide radical (CaOH)

Alec Owens ¹, Alexander Mitrushchenkov ², Sergei N. Yurchenko ¹ and Jonathan Tennyson ¹★

¹*Department of Physics and Astronomy, University College London, Gower Street, WC1E 6BT London, UK*

²*MSME, Université Gustave Eiffel, CNRS UMR 8208, Univ Paris Est Creteil, F-77474 Marne-la-Vallée, France*

Accepted 2022 August 26. Received 2022 August 25; in original form 2022 July 13

ABSTRACT

Any future detection of the calcium monohydroxide radical (CaOH) in stellar and exoplanetary atmospheres will rely on accurate molecular opacity data. Here, we present the first comprehensive molecular line list of CaOH covering the $\tilde{A}^2\Pi-\tilde{X}^2\Sigma^+$ rotation-vibration-electronic and $\tilde{X}^2\Sigma^+-\tilde{X}^2\Sigma^+$ rotation-vibration bands. The newly computed OYT6 line list contains over 24.2 billion transitions between 3.2 million energy levels with rotational excitation up to $J = 175.5$. It is applicable to temperatures up to $T = 3000$ K and covers the $0-35\,000\text{ cm}^{-1}$ range (wavelengths $\lambda > 0.29\ \mu\text{m}$) for rotational, rotation-vibration and the $\tilde{A}^2\Pi-\tilde{X}^2\Sigma^+$ electronic transitions. The strong band around $16\,000\text{ cm}^{-1}$ ($\lambda = 0.63\ \mu\text{m}$) is likely to be of interest in future astronomical observations, particularly in hot rocky exoplanets where temperatures can become extremely high. The OYT6 line list has been generated using empirically refined $\tilde{X}^2\Sigma^+$ and $\tilde{A}^2\Pi$ state potential energy surfaces, high-level *ab initio* transition dipole moment surfaces, and a rigorous treatment of both Renner–Teller and spin-orbit coupling effects, which are necessary for correctly modelling the CaOH spectrum. Post-processing of the CaOH line list has been performed so as to tailor it to high-resolution applications, i.e. by replacing calculated energy levels with more accurate empirically derived values (where available), hence improving the accuracy of the predicted line positions in certain regions. The OYT6 line list is available from the ExoMol data base at www.exomol.com and the CDS astronomical data base.

Key words: molecular data – opacity – planets and satellites: atmospheres – stars: atmospheres – ISM: molecules.

1 INTRODUCTION

The calcium monohydroxide radical ($^{40}\text{Ca}^{16}\text{O}^1\text{H}$) has been cited as a missing opacity source in the study of M-dwarf photospheres using the BT–Settl model (Rajpurohit et al. 2013). The atmospheres of hot rocky super-Earth exoplanets are also expected to contain spectroscopic signatures of CaOH (Bernath 2009; Tennyson & Yurchenko 2017) but a lack of molecular line list data in the necessary spectral regions inhibits detections. Notably, it is the rotation-vibration-electronic (rovibronic) spectrum composed of the strong $\tilde{A}^2\Pi-\tilde{X}^2\Sigma^+$ band system that is of interest. This has motivated the ExoMol data base (Tennyson & Yurchenko 2012; Tennyson et al. 2016, 2020) to work on the production of a high-accuracy line list for CaOH and this paper is the final stage in that process.

In previous work, we performed an extensive literature search and extracted all meaningful rovibronic spectroscopic transitions (around 3200) of CaOH (Wang et al. 2020) and processed them using the MARVEL (measured active rotation vibration energy levels) algorithm (Császár et al. 2007; Furtenbacher, Császár & Tennyson 2007; Furtenbacher & Császár 2012; Tóbiás et al. 2019). This procedure produced a consistent data set of empirical-quality energy levels, each possessing a measurement uncertainty and unique quantum number labelling. The advantages of having such a data

set are twofold: First, the theoretical spectroscopic model of the molecule, primarily the molecular potential energy surface (PES), can be fine-tuned in calculations so as to improve the accuracy of the predicted line positions and to a lesser extent the line intensities. Secondly, when post-processing the final line list, the computed energy levels can be substituted with the more accurate MARVEL values (where available). Doing so can dramatically improve the accuracy of certain line positions, making the final list suitable in certain windows for studying exoplanet atmospheres at high spectral resolution (Snellen 2014; Birkby 2018). Currently, the adaptation of molecular line lists in the ExoMol data base for high-resolution applications is being actively pursued, for example, see Bowesman et al. (2021).

A full-dimensional *ab initio* spectroscopic model of the $\tilde{A}^2\Pi-\tilde{X}^2\Sigma^+$ band system of CaOH has also been reported by the authors (Owens et al. 2021a). High-level *ab initio* theory was used to construct new potential energy and transition dipole moment surfaces, and both Renner–Teller and spin-orbit coupling effects, which are essential for correctly reproducing the CaOH spectrum, were accounted for in the calculations. Whilst the ground $\tilde{X}^2\Sigma^+$ PES was empirically refined to the MARVEL data set of energy levels, vastly improving its accuracy, the $\tilde{A}^2\Pi$ state PES was not rigorously refined as further development of the computer program EVEREST (Mitrushchenkov 2012) was necessary. This has now been completed for this work, ultimately leading to considerably more accurate $\tilde{A}^2\Pi$ state energy levels and $\tilde{A}-\tilde{X}$ transition wavenumbers

* E-mail: j.tennyson@ucl.ac.uk

(up to several orders-of-magnitude). Doing so has paved the way for the production of a comprehensive, high-temperature line list of CaOH which we report below.

It is worth mentioning that the spectroscopic model and associated $\tilde{A} - \tilde{X}$ line list produced in this work could greatly assist the design of laser cooling schemes in ultracold molecule research and precision tests of fundamental physics. These fields have been focusing on the alkaline Earth monohydroxide radicals, notably CaOH (Augustovičová & Bohn 2019; Kozyryev et al. 2019), because of the favourable rovibronic energy level structure. Already, direct laser cooling of CaOH to temperatures near 1 millikelvin has been performed in a one-dimensional magneto-optical trap (Baum et al. 2020) and robust experimental schemes to extend cooling of CaOH into the microkelvin regime have been suggested (Baum et al. 2021) and successfully implemented (Vilas et al. 2022). The calculation of Einstein A coefficients between the different molecular states in CaOH would allow branching ratios and decay routes to be analysed, and such work has previously been carried out for other molecules using ExoMol line lists (Owens et al. 2017).

2 METHODS

2.1 Refinement of the $\tilde{A}^2\Pi$ PES

In CaOH, the Renner–Teller effect (Renner 1934) causes the $\tilde{A}^2\Pi$ state to split into two Renner surfaces A' and A'' at bent molecular geometries. Any correct description of the spectrum of CaOH must account for the Renner–Teller effect (Hilborn, Qingshi & Harris 1983; Coxon, Li & Presunka 1991; Jarman & Bernath 1992; Li & Coxon 1992; Coxon, Li & Presunka 1994; Li & Coxon 1995, 1996). In our spectroscopic model, the A' and A'' surfaces are each represented by their own analytic function, essentially an eighth-order polynomial in terms of linear expansion variables of the bond lengths and bond angle, see Owens et al. (2021a) for a detailed description. The A' and A'' potential functions share the same parameter values for stretching only terms but possess different parameter values for terms involving the bending mode.

Despite numerous studies of the $\tilde{A} - \tilde{X}$ band system of CaOH (Bernath & Brazier 1985; Coxon et al. 1991; Coxon, Li & Presunka 1992; Li & Coxon 1992; Coxon et al. 1994; Li & Coxon 1995; Dick et al. 2006), only four vibrational levels of the $\tilde{A}^2\Pi$ state have been characterized experimentally, namely the ground vibrational state (0,0,0), the first excited Ca–O stretching ν_1 state (1,0,0), and the first and second excited bending ν_2 states (0, 1¹, 0) and (0, 2⁰, 0). Here, vibrational states are labelled with normal mode notation (ν_1, ν_2^L, ν_3), where ν_1 and ν_3 correspond to the symmetric and asymmetric stretching modes, respectively, and ν_2 labels the bending mode. The quantum number L is related to the absolute value of the vibrational angular momentum quantum number l associated with the ν_2 bending mode, $L = |l|$. It is the ν_2 bending mode which exhibits Renner–Teller splitting. No studies of the O–H stretching ν_3 mode have been carried out. As a result, we do not allow potential parameters related to the ν_3 mode to vary in the refinement and they are instead fixed to the original *ab initio* values determined in Owens et al. (2021a). Computed energy levels and transitions involving the ν_3 mode will therefore not be as accurate as those involving only the ν_1 and ν_2 modes in the final line list. This should not be overly problematic, however, as rovibronic transitions from the ground $\tilde{X}^2\Sigma^+$ electronic state to the ground, ν_1 and ν_2 vibrational levels in the $\tilde{A}^2\Pi$ state are the strongest and these levels have all been characterized by experiment and utilized in the refinement.

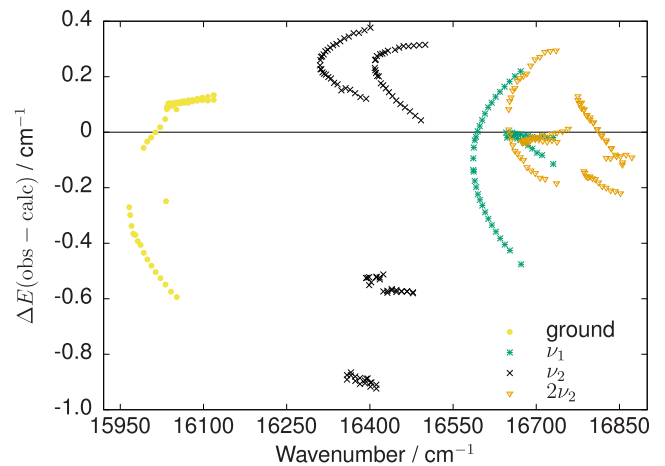


Figure 1. Residual errors $\Delta E(\text{obs} - \text{calc})$ between the empirically derived MARVEL energy levels of CaOH and the calculated EVEREST values using the refined PES for the different vibrational levels ν in the $\tilde{A}^2\Pi$ state.

The $\tilde{A}^2\Pi$ state PESs were refined using the EVEREST code and the new implementation is discussed in Section 2.2. A total of 10 potential parameters of the A' and A'' surfaces were simultaneously varied in the refinement along with the zeroth-order parameter for the spin-orbit coupling surface, which was fine-tuned to ensure accurate spin-orbit splitting in the energy level structure of CaOH. 308 term values up to $J = 15.5$, where J is the total angular momentum quantum number, were used in the refinement and were reproduced with a root-mean-square (rms) error of 0.318 cm^{-1} . A weighting scheme based on the measurement uncertainty and number of transitions involving the energy level was utilized. This information was obtained from the CaOH MARVEL analysis (Wang et al. 2020), essentially giving more importance in the refinement to well characterized term values with smaller uncertainties.

The results of the $\tilde{A}^2\Pi$ PES refinement are illustrated in Fig. 1, where we have plotted the residual errors $\Delta E(\text{obs} - \text{calc})$ between the empirically derived MARVEL and calculated energies up to $J = 15.5$. Generally speaking, the residual errors in each vibrational state exhibit similar behaviour with steadily increasing errors with increasing J (or energy as seen in Fig. 1) in both positive and negative directions. For example, in the ground vibrational state (yellow dots in Fig. 1) the residual errors for the e -parity levels become more negative with J (and energy), while the errors for the f -parity levels increase positively with J . There is also an outlier corresponding to the $J = 1.5, f$ -parity energy level at $16032.5597 \text{ cm}^{-1}$, and we suspect this to be an issue with the empirically derived MARVEL value given the consistent behaviour of the residual errors for the other energy levels in the ground vibrational state.

Interestingly, the second excited $\tilde{B}^2\Sigma^+$ state of CaOH has a minimum around 18000 cm^{-1} and couples to the $\tilde{A}^2\Pi$ state through spin-orbit coupling and linear vibronic coupling. As we have neglected the $\tilde{B}^2\Sigma^+$ state and its associated couplings in our spectroscopic model, there will no doubt be an impact on the accuracy that could be achieved had we been able to incorporate them. Even adding a ‘dummy’ coupling curve as a free fitted parameter can improve the accuracy of the spectroscopic model as is commonly seen for diatomics studied by the ExoMol project, for example, see CaH (Owens et al. 2022). The $\tilde{A}^2\Pi$ state is expected to acquire a significant fraction of $\tilde{B}^2\Sigma^+$ character through the vibronic interaction (Coxon et al. 1994) and recently it was shown that the inclusion of the $\tilde{B} - \tilde{A}$ couplings can improve the calculation of

vibronic branching ratios in CaOH (Zhang et al. 2021), which are needed when designing efficient laser-cooling schemes. That said, the effect of the $\tilde{B}^2\Sigma^+$ state on the $\tilde{A}^2\Pi$ state is to some extent absorbed into the model by empirically refining the $\tilde{A}^2\Pi$ PESs.

2.2 Extension of the EVEREST code

The EVEREST code (Mitrushchenkov 2012) was used in all rotation-vibration (rovibrational) and rovibronic line list calculations and for the $\tilde{A}^2\Pi$ PES refinement. EVEREST is capable of treating interacting electronic states with spin-dependent coupling and Renner–Teller effects in triatomic molecules. The program is general in its design and employs an exact kinetic energy operator. Extending the functionality of EVEREST to be able to perform PES refinement required several developments.

The PESs of the ground state V_X and excited states $V_{A'}$, $V_{A''}$ are represented using the Taylor-type expansion

$$V_S = \sum_{ijk} f_{ijk}^{(S)} \xi_1^i \xi_2^j \xi_3^k, \quad S = X, A', A'', \quad (1)$$

where the coordinates

$$\xi_1 = (r_1 - r_1^{\text{eq}})/r_1, \quad (2)$$

$$\xi_2 = (r_2 - r_2^{\text{eq}})/r_2, \quad (3)$$

$$\xi_3 = \alpha - \alpha_{\text{eq}}, \quad (4)$$

with the internal stretching coordinates $r_1 = r_{\text{CaO}}$, $r_2 = r_{\text{OH}}$, the interbond angle $\alpha = \angle(\text{CaOH})$, and the equilibrium parameters r_1^{eq} , r_2^{eq} and α_{eq} [see Owens et al. (2021a)].

The refinement procedure is a least-squares fitting of the potential parameters $f_{ijk}^{(S)}$ and structural parameters r_i^{eq} and α_{eq} to the observed (MARVEL) energies of CaOH, realized through the minimization of the functional

$$F = \sum_i w_i \left(E_i^{(\text{obs})} - E_i^{(\text{calc})} \right)^2, \quad (5)$$

using Newton’s steepest descent method. Here $E_i^{(\text{obs})}$ and $E_i^{(\text{calc})}$ are the observed and calculated rovibronic energies, respectively, and w_i are fitting weights. The first derivatives of $E_i \equiv E_i^{(\text{calc})}$ with respect to (wrt) the varied potential parameters $f_{ijk}^{(S)}$ (required in the least-squares fitting) are evaluated using the Hellmann–Feynman theorem such that the derivatives of the rovibronic energies can be expressed as expectation values

$$\frac{\partial E_n}{\partial f_{ijk}^{(S)}} = \langle n | \frac{\partial \hat{H}}{\partial f_{ijk}^{(S)}} | n \rangle, \quad (6)$$

where $|n\rangle$ is the corresponding rovibronic EVEREST wavefunction. This method takes advantage of the linear form of the $f_{ijk}^{(S)}$ parametrization,

$$\langle n | \frac{\partial \hat{H}}{\partial f_{ijk}^{(S)}} | n \rangle = \langle n | \xi_1^i \xi_2^j \xi_3^k | n \rangle. \quad (7)$$

Some of the A' and A'' potential parameters of the two Renner–Teller components corresponding to the linear geometry are shared between the two surfaces, namely the stretch-only terms $f_{ij0}^{(A')} = f_{ij0}^{(A'')} \equiv f_{ij0}^{(A)}$. The corresponding derivatives are given by the chain rule as follows,

$$\frac{\partial E_n}{\partial f_{ij0}^{(A)}} = \frac{\partial E_n}{\partial f_{ij0}^{(A')}} + \frac{\partial E_n}{\partial f_{ij0}^{(A'')}}. \quad (8)$$

The derivatives wrt to the non-linear parameters r_1^{eq} , r_2^{eq} and α_{eq} are obtained using finite differences, for example,

$$\frac{\partial E_n}{\partial r_1^{\text{eq}}} = \frac{E_n(r_1^{\text{eq}} + \Delta r) - E_n(r_1^{\text{eq}} - \Delta r)}{2\Delta r}, \quad (9)$$

where Δr is a small displacement, typically taken as 0.1 per cent of the parameter value, 0.002 Å or 0.2° for the bond lengths or inter-bond angles, respectively.

The implementation of the refinement procedure closely follows a previously developed Fortran wrapper for the triatomic nuclear motion program DVR3D reported in Polyansky et al. (2018), where the underlying *ab initio* PES was used to constrain the shape of the refined PES via a simultaneous fit to the original *ab initio* energies. The fitting weights were initially set to be inversely proportional to the experimental (MARVEL) uncertainties and then adjusted via Watson’s robust weighting scheme (Watson 2003). The expectation values of the potential expansions are now a part of the EVEREST program. The least-squares fitting procedure was implemented as an external, generalized fitting wrapper written in Fortran 95.

2.3 EVEREST line list calculations

The final list calculations in EVEREST used the same set up and parameters as the PES refinement. Specifically, valence bond length – bond angle coordinates with a discrete variable representation (DVR) basis composed of 100 Sinc-DVR functions on both the Ca–O bond in the 2.6–7.0 a_0 interval and the O–H bond in the 1.1–6.0 a_0 interval, and with 120 Legendre functions for the $\angle(\text{CaOH})$ bond angle. The $J = 0$ vibrational eigenfunctions were computed up to 10000 cm^{-1} above the lowest vibronic state for $0 \leq K \leq 27$, where $K = |\Lambda + l|$ (Λ and l are the projections of the electronic and vibrational angular momenta along the linear axis) from a Hamiltonian with a dimension of 10000. For $K \geq 1$, the Renner–Teller effect was explicitly taken into account by solving the coupled A'/A'' problem, see Mitrushchenkov (2012) for the full methodology. The full rovibronic Hamiltonian including spin-orbit coupling was built and diagonalized using these vibronic states for J up to 175.5, where J is the total angular momentum quantum number. Convergence of the computed rovibronic states was checked by running test calculations with $K \leq 50$, increasing the Hamiltonian dimension, and using larger DVR grids for the coordinates. The final selected EVEREST calculation parameters produced sufficiently converged results and ensured that calculations were computationally tractable. All rovibrational transitions in the $\tilde{X}^2\Sigma^+ - \tilde{X}^2\Sigma^+$ band and rovibronic transitions in the $\tilde{A}^2\Pi - \tilde{X}^2\Sigma^+$ band were calculated to produce a line list containing 24,215,753,701 transitions between 3,187,522 states with rotational excitation up to $J = 175.5$. The CaOH line list covers the 0–35 000 cm^{-1} range (wavelengths $\lambda > 0.29 \mu\text{m}$).

3 RESULTS

3.1 Line list format

The ExoMol data format has been discussed in detail (Tennyson et al. 2020) and is used for all the line lists in the ExoMol data base. The `.trans` file, see Table 1, contains the computed transitions with upper and lower state ID labels, Einstein A coefficients (in s^{-1}) and transition wavenumbers (in cm^{-1}). The `.states` file, see Table 2, contains the computed rovibronic energy levels (in cm^{-1}), each labelled with a unique state ID counting number and quantum number labelling.

Table 1. Extract from the `.trans` file of the CaOH OYT6 line list.

f	i	A_{fi}
419	1	2.8829E-08
2453	419	2.8829E-08
61925	59544	1.0353E-15
61926	59543	1.0353E-15
73928	61926	6.8168E-18

Notes. f : Upper state counting number;
 i : Lower state counting number;
 A_{fi} : Einstein- A coefficient (in s^{-1}).

To adapt the CaOH line list for high-resolution applications we have replaced the computed EVEREST energy levels and their uncertainties with the empirical MARVEL values where available. This was done for a total of 1614 states up to $J = 61.5$ for the ground $\tilde{X}^2\Sigma^+$ and excited $\tilde{A}^2\Pi$ states. The final column in the `.states` file indicates whether the energy level is either calculated (Ca) or from MARVEL (Ma). All computed energy levels have an estimated uncertainty of 10 cm^{-1} , which is a conservative estimate to ensure that users are aware of the difference in reliability between calculated and ‘MARVELised’ lines, particularly for high-resolution applications. The penultimate column of the `.states` file is the EVEREST calculated energy for reference. The quantum numbers used to label the rovibronic states of CaOH are listed in Table 2 and discussed in detail in Wang et al. (2020).

3.2 Temperature-dependent partition functions

Intensity simulations require knowledge of the temperature-dependent partition function $Q(T)$, defined as

$$Q(T) = \sum_i g_i \exp\left(\frac{-E_i}{kT}\right), \quad (10)$$

where $g_i = g_{\text{ns}}(2J_i + 1)$ is the degeneracy of a state i with energy E_i and total angular momentum quantum number J_i , and the nuclear spin statistical weight $g_{\text{ns}} = 2$ for CaOH. Partition function values were computed on a 1 K grid in the 1–4000 K range and can be downloaded from the ExoMol website at www.exomol.com along with the OYT6 line list. The convergence of $Q(T)$ as a function of J is shown in Fig. 2 for different temperatures. At $T = 3000\text{ K}$, the convergence of the partition function with J is slower and we recommend this as a soft upper limit for using the OYT6 line list at elevated temperatures. Application of the OYT6 line list above this temperature may result in a progressive loss of opacity.

In Fig. 3, we have compared partition function values up to $T = 4000\text{ K}$ computed using the OYT6 line list against the JANAF CaOH values (Irwin 1988). The JANAF values were calculated using a simple molecular model that employs a rigid rotor, simple harmonic oscillator approximation. As such, although the general behaviour between the OYT6 and JANAF values is in reasonable agreement, we expect the OYT6 values to be more accurate at least up to $T = 3000\text{ K}$. This is because the OYT6 values are computed by summing over all meaningful rotation-vibration energy levels in the lowest two electronic states of CaOH. Interestingly, the contribution from $\tilde{A}^2\Pi$ state energy levels in the summation of equation (10) to the partition

function is near-negligible. For example, the partition function at $T = 3000\text{ K}$ calculated using only ground $\tilde{X}^2\Sigma^+$ state energy levels in the summation of equation (10) showed a 0.08 per cent difference to the full OYT6 value and this is the largest difference encountered.

3.3 Simulated spectra

There is no meaningful intensity information in the literature on the $\tilde{A} - \tilde{X}$ band system that we can compare to, making it difficult to evaluate the accuracy of our computed line intensities against measured spectra. Transition intensities computed using *ab initio* dipole moment surfaces (DMSs) are comparable to, and in certain instances, more reliable than experiment (Tennyson 2014; Yurchenko 2014). As an estimate we would expect the CaOH line intensities to be within 5–10 per cent of experimentally determined values. It is also worth noting that strong electronic transitions are less sensitive to the shape of the transition DMSs, given that vibronic intensities can be relatively well modelled using only a single value for the transition dipole times a Franck–Condon factor.

In Fig. 4, the temperature dependence of the CaOH spectrum is illustrated, and we have simulated absolute absorption cross-sections at a resolution of 1 cm^{-1} using a Gaussian line profile with a half width at half maximum (HWHM) of 1 cm^{-1} . All spectral simulations were done using the EXOCROSS program (Yurchenko, Al-Refaie & Tennyson 2018b). The strongest rovibronic features of the OYT6 line list are around $16\,000\text{ cm}^{-1}$ and even at high temperatures ($T = 3000\text{ K}$) these still dominate the CaOH spectrum. This region is significant for atmospheric studies of hot rocky super-Earths, a class of exoplanets that are in close proximity to their host star and exposed to extremely high temperatures, e.g. up to 4000 K .

Closer inspection of the main infrared part of the spectrum is shown in Fig. 5 at $T = 1000\text{ K}$. The fundamental wavenumbers of the $\tilde{X}^2\Sigma^+$ state of CaOH occur at $\nu_1 \approx 609$, $\nu_2 \approx 353$, and $\nu_3 \approx 3792\text{ cm}^{-1}$ (Owens et al. 2021a). The spectrum exhibits a relatively slow flattening of band intensities with increasing wavenumber and similar behaviour was also observed in line list calculations of the alkali metal hydroxides KOH and NaOH [Owens, Tennyson & Yurchenko (2021b)], which both possess linear molecular structures.

In Fig. 6, absolute absorption line intensities (in units of $\text{cm}/\text{molecule}$) at $T = 700\text{ K}$ have been plotted for the strongest regions of the OYT6 spectrum. Line intensities were simulated using the expression,

$$I(f \leftarrow i) = \frac{A_{fi}}{8\pi c} g_{\text{ns}}(2J_f + 1) \frac{\exp(-E_i/kT)}{Q(T)v_{fi}^2} \times \left[1 - \exp\left(-\frac{hcv_{fi}}{kT}\right) \right], \quad (11)$$

where A_{fi} is the Einstein A coefficient of a transition with wavenumber ν_{fi} (in cm^{-1}) between an initial state with energy E_i and a final state with rotational quantum number J_f . Here, k is the Boltzmann constant, h is the Planck constant, c is the speed of light and T is the absolute temperature. The nuclear spin statistical weight $g_{\text{ns}} = 2$ for CaOH and $Q(T)$ is the temperature-dependent partition function. Interestingly, the structure of the bands using the refined $\tilde{A}^2\Pi$ state PES are quantitatively different to those predicted using the original *ab initio* PES, see Fig. 5 in Owens et al. (2021a). This behaviour has been encountered before by the ExoMol project and highlights the fact that rovibronic spectra can be very sensitive to the quality of the underlying PESs with both line positions and intensities being affected.

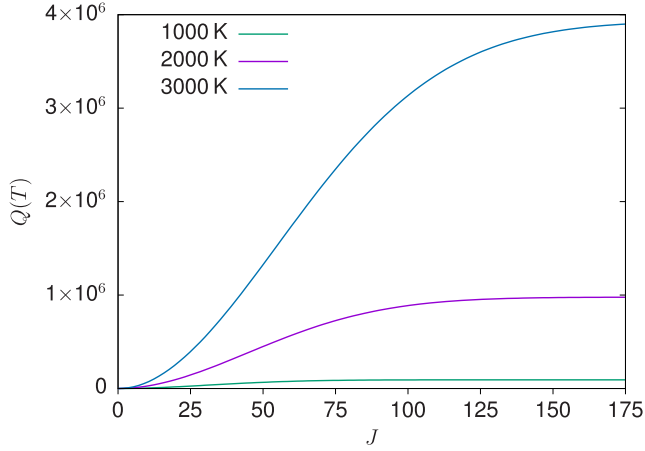


Figure 2. Convergence of the partition function $Q(T)$ of CaOH wrt the total angular momentum quantum number J at different temperatures.

In Fig. 7, a CaOH OYT6 stick spectrum at $T = 300$ K, computed using equation (11), is plotted against pure rotational transitions in the ground $\tilde{X}^2\Sigma^+$ state extracted from the Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2001; Müller et al. 2005). The rotational band is reproduced well with the OYT6 spectrum possessing more structure and lines since it extends up to $J = 175.5$, compared to CDMS which contains 102 lines up to $J = 50.5$. The CDMS line intensities are noticeably stronger, however, and they were computed using a more simplified model based on a single dipole moment value of $\mu = 1.47$ Debye (Steimle et al. 1992). Our value for the ground $\tilde{X}^2\Sigma^+$ state dipole is $\mu = 1.09$ Debye and this was computed using a high-level *ab initio* theory (restricted coupled cluster with a large augmented correlation consistent basis set RCCSD(T)/aug-pwCVQZ-PP). This value is closer to other CaOH dipole values of 0.98 Debye from *ab initio*

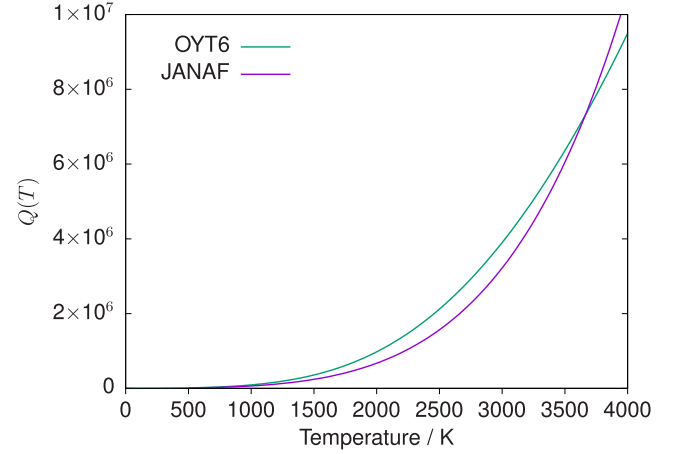


Figure 3. Computed OYT6 partition function values $Q(T)$ of CaOH for temperatures up to $T = 4000$ K compared against the JANAF values (Irwin 1988).

calculations (Bauschlicher et al. 1990) and a semi-empirical value of 1.2 Debye (Mestdagh & Visticot 1991), and would explain the weaker OYT6 line intensities. These comparisons also give us confidence in the reliability of the DMSs used in intensity calculations for the OYT6 line list.

4 CONCLUSION

A new comprehensive molecular line list of the $\tilde{X}^2\Sigma^+ - \tilde{X}^2\Sigma^+$ rovibrational and $\tilde{A}^2\Pi - \tilde{X}^2\Sigma^+$ rovibronic band system of CaOH has been reported. The OYT6 line list covers rotational, rotation-vibration, and the $\tilde{A}^2\Pi - \tilde{X}^2\Sigma^+$ electronic transition in the 0–35 000 cm^{-1} range (wavelengths $\lambda > 0.29 \mu\text{m}$) and is applicable to temperatures up to $T = 3000$ K. For CaOH, we expect the

Table 2. Extract from the .states file of the CaOH OYT6 line list. For the description of the quantum numbers see Wang et al. (2020).

i	\tilde{E} (cm^{-1})	g_i	J	unc	τ	elf	N	State	L	v_1	v_2	l_2	v_3	Ω	F_i	Calc.	Lab.
4056	16032.559740	8	1.5	0.014003	−1	e	1	A'	1	0	0	0	0	1.5	F2	16032.218598	Ma
4057	16311.011570	8	1.5	0.007071	−1	e	1	A'	0	0	1	1	0	0.5	F1	16311.301542	Ma
4058	16325.468962	8	1.5	10.000000	−1	e	2	A'	2	0	1	1	0	1.5	F1	16325.468962	Ca
4059	16410.656140	8	1.5	0.010000	−1	e	2	A'	0	0	1	1	0	0.5	F1	16410.856088	Ma
4060	16587.366240	8	1.5	0.005000	−1	e	2	A'	1	1	0	0	0	0.5	F1	16587.401835	Ma
4061	16645.654491	8	1.5	10.000000	−1	e	1	A'	1	1	0	0	0	1.5	F2	16645.654491	Ca
4062	16650.410240	8	1.5	0.010000	−1	e	2	A'	1	0	2	0	0	0.5	F1	16650.559911	Ma

Notes: i : State counting number.

\tilde{E} : State energy (in cm^{-1}).

g_i : Total statistical weight, equal to $g_{ns}(2J + 1)$.

J : Total angular momentum.

unc: Uncertainty (in cm^{-1}).

τ : Total parity.

N : Rotational angular momentum.

State: Electronic state X, A', A'' .

L : Vibronic angular momentum quantum number, used to label vibronic states with $\Sigma, \Pi, \Delta, \Phi, \dots$ for $L = 0, 1, 2, 3, \dots$

$L = l_{\text{vib}} + \Lambda$, where l_{vib} and Λ are projections of the vibrational and electronic angular momenta on the z axis.

v_1 : Symmetric stretching v_1 mode vibrational quantum number ('−1' means unassigned).

v_2 : Bending v_2 mode vibrational quantum number ('−1' means unassigned).

l_2 : Vibrational angular momentum quantum number associated with v_2 mode ('−1' means unassigned).

v_3 : Antisymmetric stretching v_3 mode vibrational quantum number ('−1' means unassigned).

Ω Projection of the total angular momentum $|\Omega| \leq J$.

F_i : Spin components F_1 and F_3 .

Calc: Original EVEREST calculated state energy (in cm^{-1}).

Label: Label 'Ma' for MARVEL, 'Ca' for calculated.

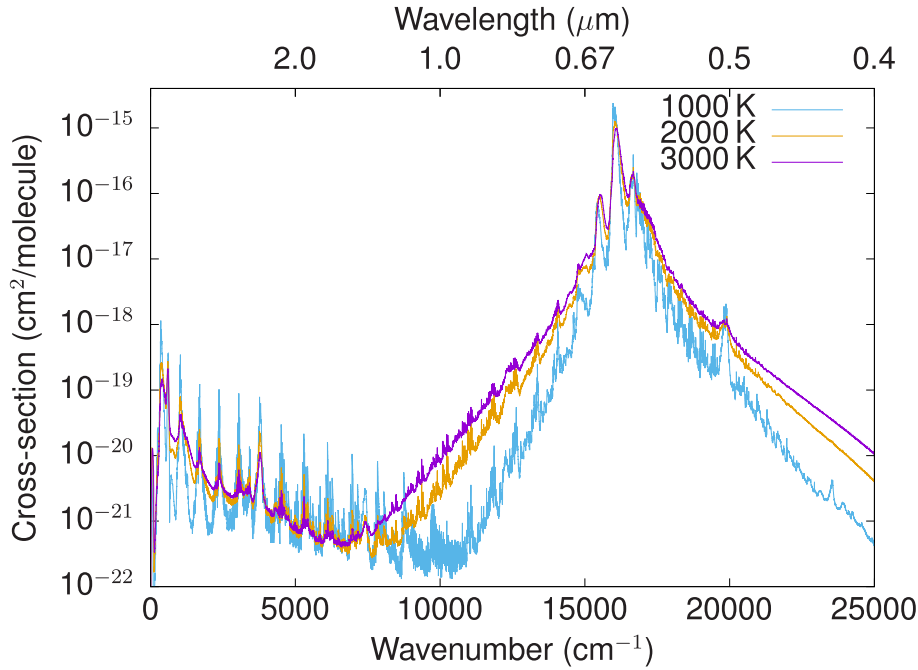


Figure 4. Spectrum of CaOH at various temperatures (in K). Absorption cross-sections were computed at a resolution of 1 cm^{-1} and modelled with a Gaussian line profile with a HWHM of 1 cm^{-1} .

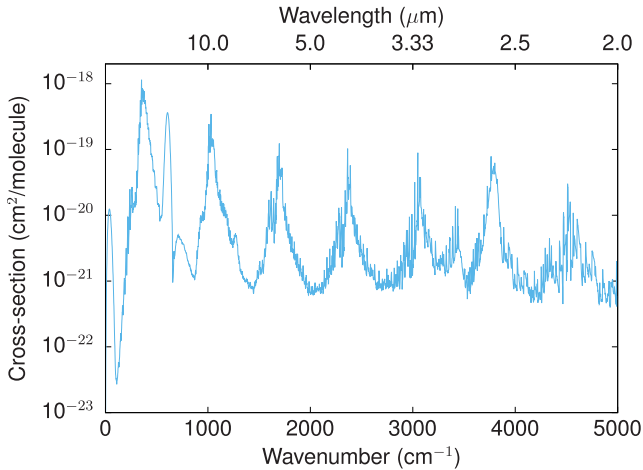


Figure 5. Absorption cross-sections at a temperature of $T = 1000 \text{ K}$ in the $0\text{--}5000 \text{ cm}^{-1}$ range at a resolution of 1 cm^{-1} , modelled with a Gaussian line profile with a HWHM of 1 cm^{-1} .

strong $\tilde{A}^2\Pi\text{--}\tilde{X}^2\Sigma^+$ band around 16000 cm^{-1} ($\lambda = 0.63 \mu\text{m}$) to be of significant interest as it dominates the spectrum at very high temperatures, which will be relevant to the conditions found on hot rocky super-Earth exoplanets. Also important is that the OYT6 line list has been adapted to high-resolution applications by replacing computed energy levels with the more accurate empirically derived MARVEL values where available. To generate the OYT6 line list, extension of the computer program EVEREST was necessary such that the $\tilde{A}^2\Pi$ state PESs of CaOH could be rigorously refined to experimental data, thus considerably improving the accuracy of the computed spectra. The extension of EVEREST will also lend itself to future spectroscopic investigations, for example, the near-ultraviolet (UV) spectrum of water which is of considerable interest due to

disagreement over absorption in this region, see Wang et al. (2022) and references within. Any future line list of H_2O in the near-UV region would find great applicability in exoplanetary research.

Rajpurohit et al. (2013) used BT-Settl (Allard 2014) to analyse synthetic spectra of M-dwarf stars; they suggested that BT-Settl was missing key absorption due to three molecules: AlH and NaH in blue, and CaOH at 5570 \AA . Previous ExoMol studies have provided line lists and opacities for the metal hydrides AlH (Yurchenko et al. 2018a) and NaH (Rivlin et al. 2015); provision of the OYT6 line list for CaOH means that the set should now be complete.

It is worth mentioning that the lowest-lying eight electronic states of CaOH up to the $\tilde{G}^2\Pi$ state (Hailey, Jarman & Bernath 1997) at approximately 32633 cm^{-1} are known. While it is unclear how strong these electronic transitions are in relation to the $\tilde{A}^2\Pi\text{--}\tilde{X}^2\Sigma^+$ band, the $\tilde{B}^2\Sigma^+\text{--}\tilde{X}^2\Sigma^+$ band would almost certainly be observable as it lies in the same spectroscopic region (approximately 18000 cm^{-1}). Future work by the ExoMol data base could address this electronic transition if deemed worthwhile.

ACKNOWLEDGEMENTS

This work was supported by the STFC Projects No. ST/M001334/1 and ST/R000476/1. The authors acknowledge the use of the UCL Legion High Performance Computing Facility (Legion@UCL) and associated support services in the completion of this work, along with the Cambridge Service for Data Driven Discovery (CSD3), part of which is operated by the University of Cambridge Research Computing on behalf of the STFC DiRAC HPC Facility (www.dirac.ac.uk). The DiRAC component of CSD3 was funded by BEIS capital funding via STFC capital grants ST/P002307/1 and ST/R002452/1 and STFC operations grant ST/R00689X/1. DiRAC is part of the National e-Infrastructure. This work was also supported by the European Research Council (ERC) under the European Union’s Horizon 2020 research and innovation programme through Advance Grant number 883830.

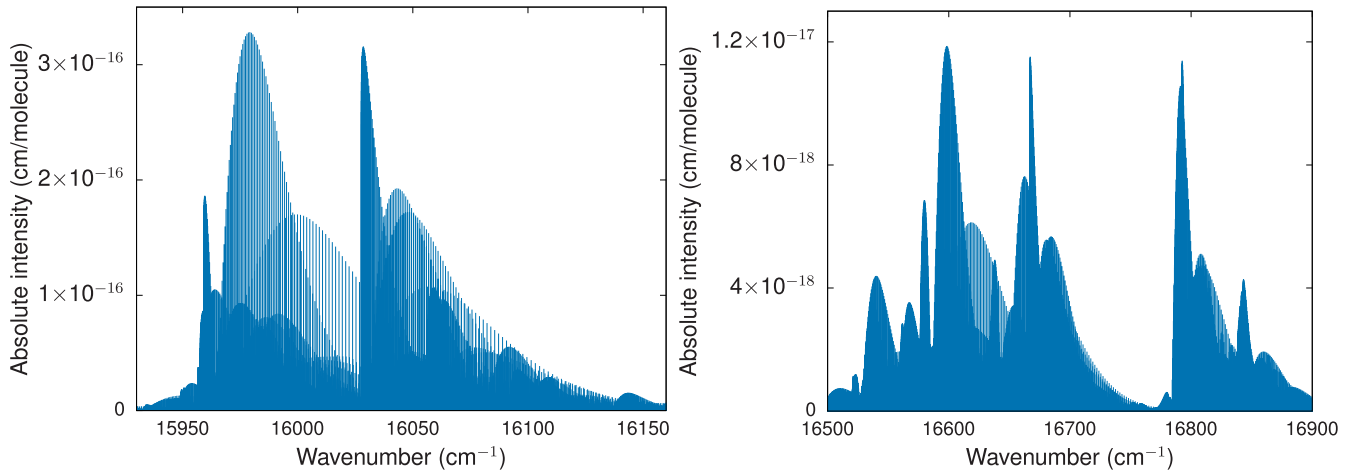


Figure 6. Absolute line intensities of the strongest $\tilde{A}^2\Pi$ bands of CaOH at a temperature of $T = 700$ K.

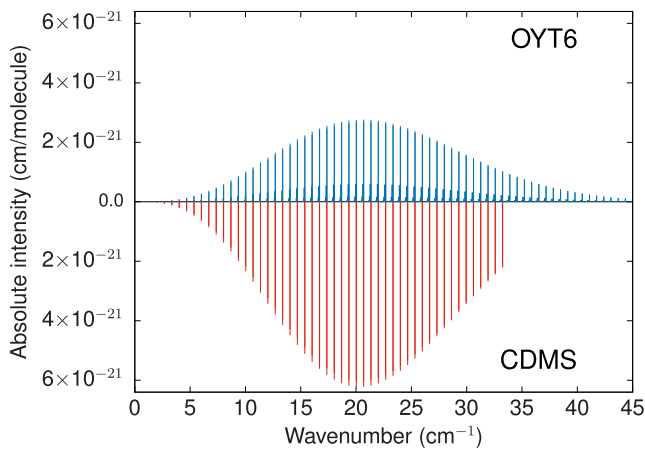


Figure 7. Comparison of the OYT6 stick spectrum at $T = 300$ K in the microwave region against all transition data from the CDMS (Müller et al. 2001; Müller et al. 2005).

DATA AVAILABILITY

The states, transition and partition function files for the OYT6 line list can be downloaded from www.exomol.com and the CDS data centre cds.unistra.fr. The open access program EXOCROSS is available from github.com/exomol.

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