

Water vapour in the atmosphere of the habitable-zone eight Earth-mass planet K2-18 b

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1 In the past decade, observations from space and ground have found H₂O to be the most
2 abundant molecular species, after hydrogen, in the atmospheres of hot, gaseous, extrasolar
3 planets¹⁻⁵. Being the main molecular carrier of oxygen, H₂O is a tracer of the origin and
4 the evolution mechanisms of planets. For temperate, terrestrial planets, the presence of H₂O
5 is of great significance as an indicator of habitable conditions. Being small and relatively
6 cold, these planets and their atmospheres are the most challenging to observe, and there-
7 fore no atmospheric spectral signatures have so far been detected⁶. Super-Earths – planets
8 lighter than ten M_⊕ – around later-type stars may provide our first opportunity to study
9 spectroscopically the characteristics of such planets, as they are best suited for transit obser-
10 vations. Here we report the detection of an H₂O spectroscopic signature in the atmosphere
11 of K2-18 b – an eight M_⊕ planet in the habitable-zone of an M-dwarf⁷ – with high statistical
12 confidence (ADI⁵ = 5.0, ~3.6σ^{8,9}). In addition, the derived mean molecular weight suggests
13 an atmosphere still containing some hydrogen. The observations were recorded with the
14 Hubble Space Telescope/WFC3 camera, and analysed with our dedicated, publicly avail-

15 able, algorithms^{5,9}. While the suitability of M-dwarfs to host habitable worlds is still under
16 discussion¹⁰⁻¹³, K2-18 b offers an unprecedented opportunity to get insight into the composi-
17 tion and climate of habitable-zone planets.

18 Atmospheric characterisation of super-Earths is currently within reach of the Wide Field
19 Camera 3 (WFC3) onboard the *Hubble Space Telescope* (HST), combined with the recently im-
20 plemented spatial scanning observational strategy¹⁴. The spectra of three hot transiting planets
21 with radii less than $3.0 R_{\oplus}$ have been published so far: GJ-1214 b¹⁵, HD 97658 b¹⁶ and 55 Cnc e¹⁷.
22 The first two do not show any evident transit depth modulation with wavelength, suggesting an at-
23 mosphere covered by thick clouds or made of molecular species heavier than hydrogen, while only
24 the spectrum of 55 Cnc e has revealed a light-weighted atmosphere, suggesting H/He still being
25 present. In addition, transit observations of six temperate Earth-size planets around the ultra-cool
26 dwarf TRAPPIST-1 – planets b, c, d, e, f⁶, and g¹⁸ – have not shown any molecular signatures and
27 have excluded the presence of cloud-free, H/He atmospheres around them.

28 K2-18 b was discovered in 2015 by the *Kepler* spacecraft⁷, and it is orbiting around an
29 M2.5 ($[\text{Fe}/\text{H}] = 0.123 \pm 0.157$ dex, $T_{\text{eff}} = 3457 \pm 39$ K, $M_* = 0.359 \pm 0.047 M_{\odot}$, $R_* =$
30 $0.411 \pm 0.038 R_{\odot}$)¹⁹ dwarf star, 34 pc away from the Earth. The star-planet distance of 0.1429 AU¹⁹
31 suggests a planet within the star’s habitable zone ($\sim 0.12 - 0.25$ AU)²⁰, with effective temperature
32 between 200 K and 320 K, depending on the albedo and the emissivity of its surface and/or its at-
33 mosphere. This crude estimate accounts for neither possible tidal energy sources²¹ nor atmospheric
34 heat redistribution^{11,13}, which might be relevant for this planet. Measurements of the mass and the

35 radius of K2-18 b ($M_p = 7.96 \pm 1.91 M_\oplus^{22}$, $R_p = 2.279 \pm 0.0026 R_\oplus^{19}$) yield a bulk density
36 of $3.3 \pm 1.2 \text{ g/cm}^{22}$, suggesting either a silicate planet with an extended atmosphere around or an
37 interior composition with an H_2O mass fraction lower than 50%²²⁻²⁴.

38 We analyse here eight transits of K2-18 b obtained with the WFC3 camera onboard the Hub-
39 ble Space Telescope. We used our specialised, publicly available, tools^{5,9} to perform the end-to-
40 end analysis from the raw HST data to the atmospheric parameters. The accuracy of the techniques
41 used here have been demonstrated through the largest consistently analysed catalogue of exoplan-
42 etary spectra from WFC3⁵. Details on the data analysis can be found in the Methods section. Also,
43 links to the data and the codes used can be found in the Data availability and Code availability
44 sections, respectively. Alongside with the data we provide descriptions of the data structures and
45 instructions on how to reproduce the results presented here. Our analysis resulted in the detection
46 of an atmosphere around K2-18 b with an ADI⁵ (a positively defined logarithmic Bayes Factor)
47 of 5.0, or approximately 3.6σ confidence^{8,9}, making K2-18 b the first habitable-zone planet in the
48 super-Earth mass regime ($1-10 M_\oplus$) with an observed atmosphere around it.

49 More specifically, nine transits of K2-18 b were observed as part of the HST proposals 13665
50 and 14682 (PI: Björn Benneke) and the data are available through the MAST Archive (see Data
51 Availability section). Each transit was observed during five HST orbits, with the G141 infrared
52 grism ($1.1 - 1.7 \mu\text{m}$), and each exposure was the result of 16 up-the-ramp samples in the spatially
53 scanning mode. The ninth transit observation suffered from pointing instabilities and therefore we
54 decided not to include it in this analysis. We extracted the white and the spectral light curves from

55 the reduced images, following our dedicated methodology^{5,17,25}, which has been integrated into an
56 automated, self-consistent, and user-friendly Python package named *Iraclis* (see Code Availability
57 section). No systematic variations of the white light-curve R_p/R_* appeared between the eight
58 different observations. This level of stability among the extracted broad-band transit depths is not
59 always guaranteed, as consistency problems among different observations emerged in previous
60 analyses^{5,16}.

61 In our analysis, we found that the measured mid-transit times were not consistent with
62 the expected ephemeris¹⁹. We used these results to refine the ephemeris of K2-18 b to be: $P =$
63 32.94007 ± 0.00003 days and $T_0 = 2457363.2109 \pm 0.0004$ BJD_{TDB}²⁶, where P is the period,
64 T_0 is the mid-time of the transit. However, the ephemeris calculated only from the HST data
65 is not consistent with the original detection of K2-18 b. One possibility is that the very sparse
66 data from K2 are not sufficient to give a confident result. Another possibility is that we observe
67 significant transit time variations (TTVs) caused by the other planet in the system, K2-18 c²²,
68 but more observations over a long period of time are necessary to disentangle the two scenarios.
69 In addition, we used the detrended and time-aligned – i.e. with TTVs removed – white light
70 curves to also refine the orbital parameters and found them to be: $a/R_* = 81.3 \pm 1.5$ and $i =$
71 89.56 ± 0.02 deg, where a/R_* is the orbital semi-major axis normalised to the stellar radius, and i
72 is the orbital inclination.

73 We extracted eight transmission spectra of K2-18 b and combined them, using a weighted
74 average, to produce the final spectrum (Table 1). We interpreted the planetary spectrum using

75 our spectral retrieval algorithm TauREx^{9,27} (see Code Availability section) which combines highly
76 accurate line lists²⁸ (see Data Availability section) and Bayesian analysis. At an initial stage, we
77 modelled the atmosphere of K2-18 b including all potential absorbers in the observed wavelength
78 range – i.e. H₂O, CO, CO₂, CH₄ and NH₃. However, we found that only the spectroscopic signa-
79 ture of water vapour is detected with high-confidence, so we continued our analysis only with this
80 molecule as trace-gas. We modelled the atmosphere following three approaches:

- 81 • a cloud-free atmosphere containing only H₂O and H₂/He
- 82 • a cloud-free atmosphere containing H₂O, H₂/He and N₂ (N₂ acted as proxy for “invisible”
83 molecules not detectable in the WFC3 bandpass but contributing to the mean molecular
84 weight), and
- 85 • a cloudy (flat-line model) atmosphere containing only H₂O and H₂/He.

86 We retrieved a statistically significant atmosphere around K2-18 b in all simulations (Figure
87 3), and assessed the strength of the detection using the Atmospheric Detectability Index⁵ (ADI),
88 which represents the positively defined logarithmic Bayes Factor, where the null hypothesis is a
89 model that contains no active trace gases, Rayleigh scattering or collision induced absorption – i.e.
90 a flat spectrum. The retrieval simulations yield an atmospheric detection with an ADI of 5.0, 4.7
91 and 4.0, respectively. Such ADIs correspond to approximately a 3.6, 3.5, and 3.3 σ detection^{8,9},
92 respectively. This marks the first atmosphere detected around a habitable-zone super-Earth with
93 such a high level of confidence. While the H₂O + H₂/He case appears to be the most favourable,

94 this preference is not statistically significant.

95 As far as the composition is concerned, retrieval models confirm the presence of water vapour
96 in the atmosphere of K2-18 b in all the cases studied with high statistical significance. However, it
97 is not possible to constrain either its abundance or the mean molecular weight of the atmosphere.
98 For the H₂O + H₂/He case, we found the abundance of H₂O to be between 50% and 20%, while
99 for the other two cases between 0.01% and 12.5%. The atmospheric mean molecular weight can
100 be between 5.8 and 11.5 amu in the H₂O + H₂/He case, and between 2.3 and 7.8 amu for the other
101 cases. These results indicate that a non-negligible fraction of the atmosphere is still made of H/He.
102 Additional trace-gases – e.g. CH₄, NH₃ – cannot be excluded, despite not being identified with the
103 current observations: the limiting S/N and wavelength coverage of HST/WFC3 do not allow the
104 detection of other molecules.

105 The results presented here confirm the existence of a detectable atmosphere around K2-18 b,
106 making it one of the most interesting known targets for further atmospheric characterisation with
107 future observatories, like the *James Webb Space Telescope* (0.6 μm and 28 μm) and the European
108 Space Agency ARIEL mission²⁹ (0.5 μm and 7.8 μm). The wider wavelength coverage of these
109 instruments will provide information on the presence of additional molecular species and on the
110 temperature-pressure profile of the planet, towards studying the planetary climate and potential
111 habitability. While the subject of habitability for temperate planets around late-type stars is a
112 subject of active discussion^{10–13} and real progress requires significantly improved observational
113 constraints, the analysis presented here provides the first direct observation of a molecular signature

114 from a habitable-zone exoplanet, connecting these theoretical studies to observations.

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182 **Author Contributions** A.T. performed the data analysis and developed the HST analysis soft-
183 ware *Iraclis*; I.P.W developed the atmospheric retrieval software TauREx; G.T. contributed to the
184 interpretation of the results; J.T. and S.N.Y. coordinate the ExoMol project. All authors discussed
185 the results and commented on the manuscript.

186 **Financial Interests** The authors declare no competing financial interests.

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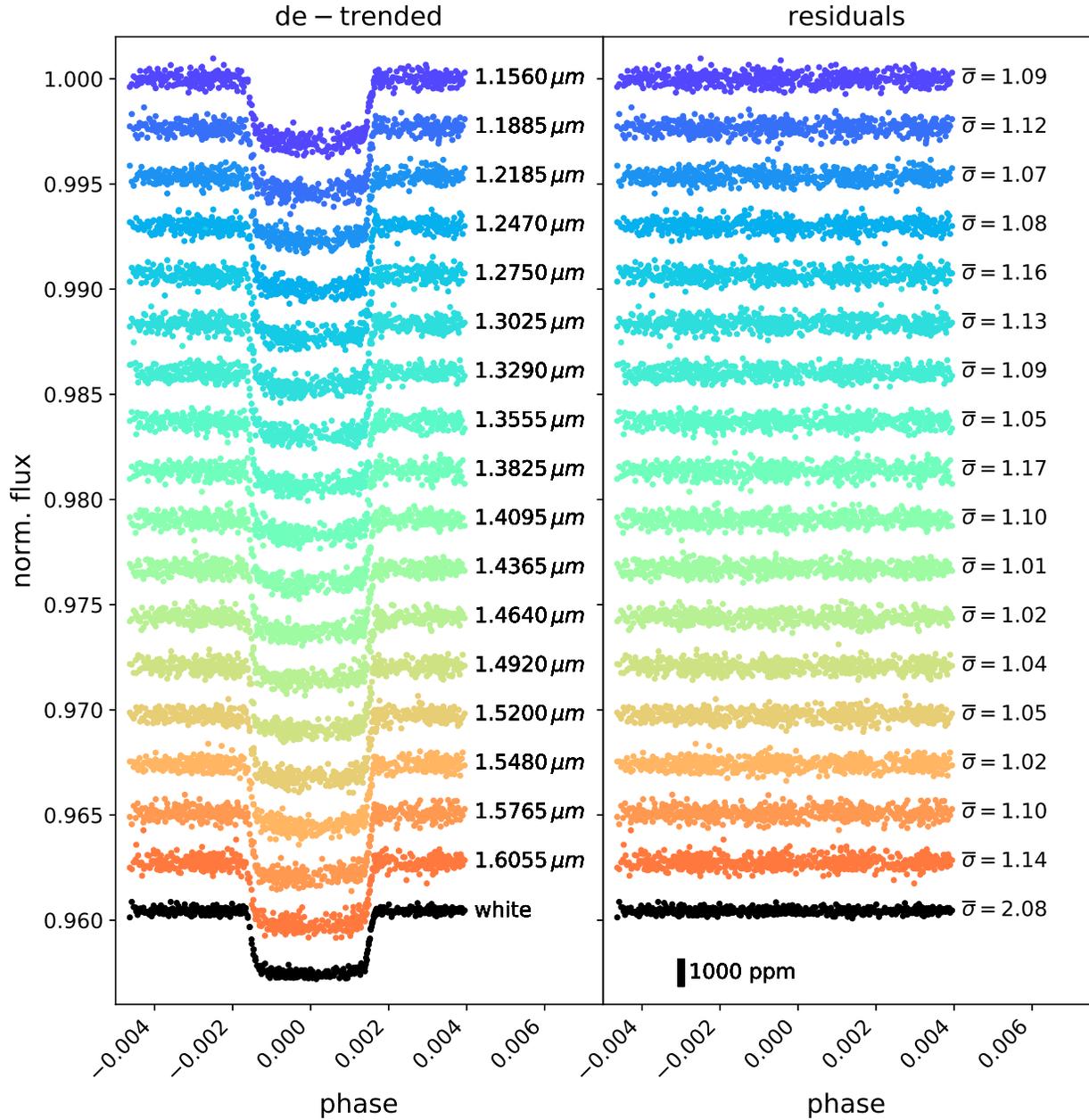


Figure 1: Analysis of the K2-18 b white (black points) and spectral (coloured points) light curves, plotted with an offset for clarity. Left: Overplotted detrended light curves. Right: Overplotted fitted residuals, where $\bar{\sigma}$ indicates the ratio between the standard deviation of the residuals and the photon noise (see Methods for more details). The black vertical bar indicates the 1000ppm scatter level.

Table 1: Transit depth $((R_p/R_*)^2)$ for the different wavelength channels, where R_p is the planetary radius, R_* is the stellar radius, and λ_1 , λ_2 are the lower and upper edges of each wavelength channel, respectively.

λ_1	λ_2	$(R_p/R_*)^2$
μm	μm	ppm
1.1390	1.1730	2905 ± 25
1.1730	1.2040	2939 ± 26
1.2040	1.2330	2903 ± 24
1.2330	1.2610	2922 ± 25
1.2610	1.2890	2891 ± 26
1.2890	1.3160	2918 ± 26
1.3160	1.3420	2919 ± 24
1.3420	1.3690	2965 ± 24
1.3690	1.3960	2955 ± 27
1.3960	1.4230	2976 ± 25
1.4230	1.4500	2990 ± 24
1.4500	1.4780	2895 ± 23
1.4780	1.5060	2930 ± 23
1.5060	1.5340	2921 ± 24
1.5340	1.5620	2875 ± 24
1.5620	1.5910	2927 ± 25
1.5910	1.6200	2925 ± 24

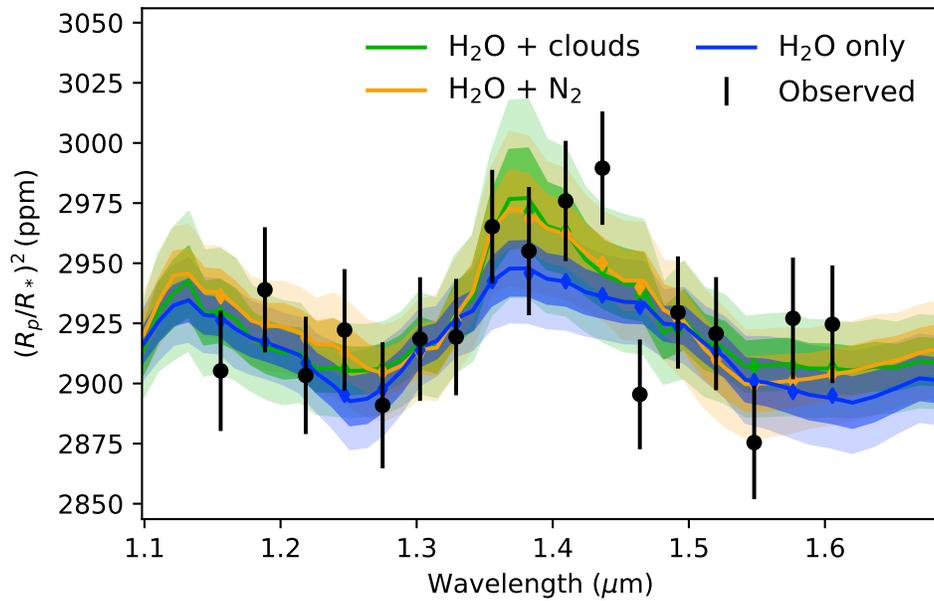
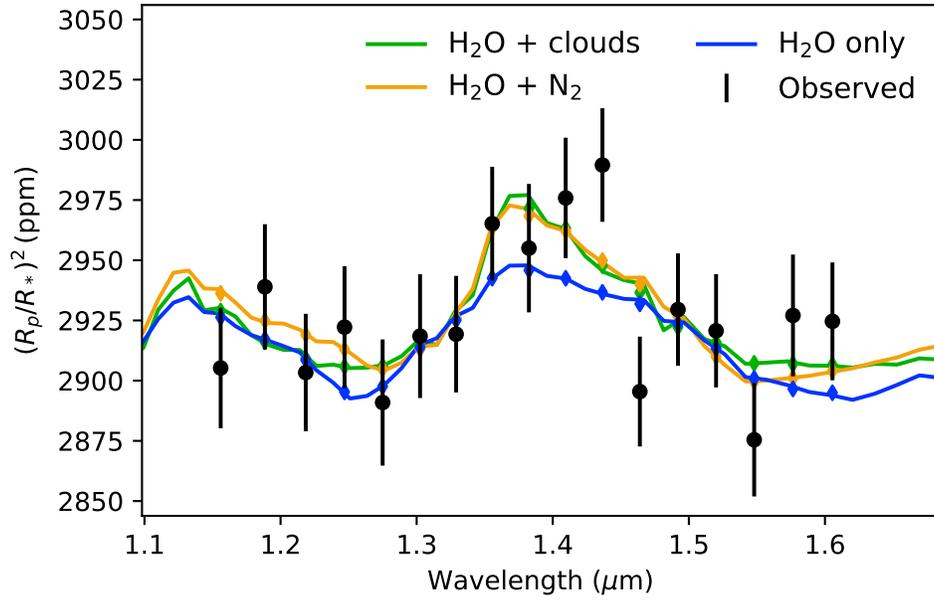


Figure 2: Best-fit models for the three different scenarios tested here: a cloud-free atmosphere containing only H₂O and H₂/He (blue), a cloud-free atmosphere containing H₂O, H₂/He and N₂ (orange), and a cloudy atmosphere containing only H₂O and H₂/He. (green). Top: best-fit models only. Bottom: 1 σ and 2 σ uncertainty ranges.

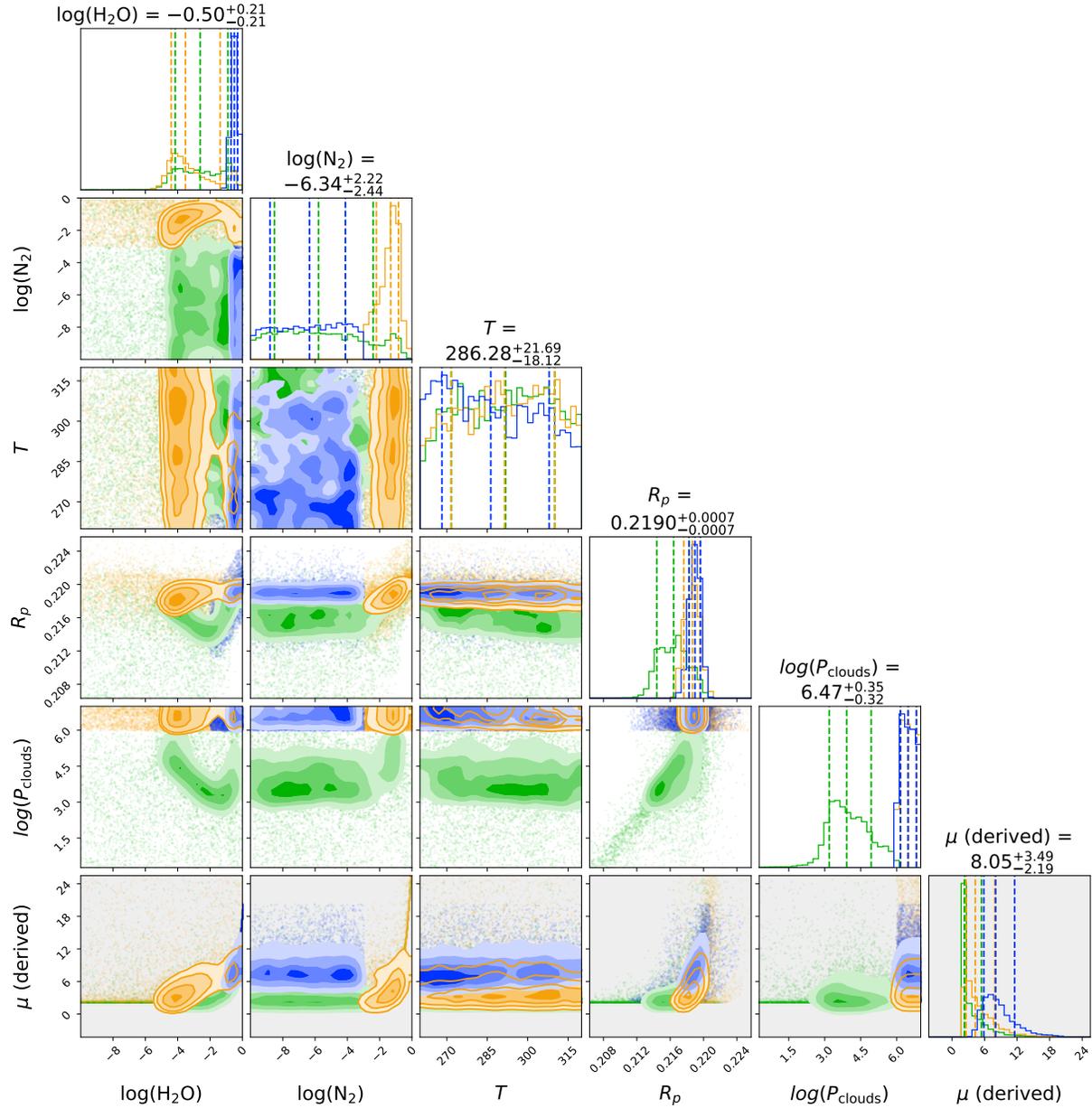


Figure 3: Posterior distributions for the three different scenarios tested here: a cloud-free atmosphere containing only H_2O and H_2/He (blue), a cloud-free atmosphere containing H_2O , H_2/He and N_2 (orange), and a cloudy atmosphere containing only H_2O and H_2/He . (green). The parameters shown, from top to bottom are: the volume mixing ratio of H_2O , the volume mixing ratio of N_2 , the planetary temperature in K, the planetary radius in R_{Jup} , the cloud top pressure in Pa, and the derived mean molecular weight.

189 **Methods**

190 **Observations** Nine transits of K2-18 b were observed as part of the HST proposals 13665 and
191 14682 (PI: Björn Benneke) and the data are available through the MAST Archive. More specif-
192 ically, the relevant HST visits are: visit 29 (06/12/2015), visit 35 (14/03/2016), and visit 30
193 (19/05/1016) from proposal 13665; visit 3 (02/12/2016), visit 1 (04/01/2017), visit 2 (06/02/2017),
194 visit 4 (13/04/2017), visit 5 (30/11/2017), and visit 6 (13/05/2018) from proposal 14682. Out of
195 these nine visits, we decided not to include the last one, as it suffered from pointing instabilities.

196 Each transit was observed during five HST orbits, with the G141 infrared grism of the WFC3
197 camera (1.1 - 1.7 μm), in the spatially scanning mode. During an exposure using the spatial scan-
198 ning mode the instrument slews along the cross-dispersion direction, allowing for longer exposure
199 times and increased signal-to-noise ratio (S/N), without the risk of saturation ¹⁴. Both forward
200 (increasing row number) and reverse (decreasing row number) scanning were used for these obser-
201 vations.

202 The detector settings were: SUBTYPE=SQ256SUB, SAMP_SEQ=SPARS10, NSAMP=16,
203 APERTURE=GRISM256, and the scanning speed was 1.4''s⁻¹. The final images had a total expo-
204 sure time of 103.128586 seconds, a maximum signal level of 1.9×10^4 electrons per pixel, and a to-
205 tal scanning length of approximately 120 pixels. Finally, for calibration reasons, a 0.833445 s non-
206 dispersed (direct) image of the target was taken at the beginning of each visit, using the F130N filter
207 and the following settings: SUBTYPE=SQ256SUB, SAMP_SEQ=RAPID, NSAMP=4, APER-

208 TURE=IRSUB256.

209 **Extracting the planetary spectrum** We carried out the analysis of the eight K2-18 b tran-
210 sits using our specialised software for the analysis of WFC3, spatially scanned spectroscopic
211 images^{5,17,25}, which has been integrated into the *Iraclis* package (see Code Availability section).
212 The reduction process included the following steps: zero-read subtraction, reference pixels correc-
213 tion, nonlinearity correction, dark current subtraction, gain conversion, sky background subtrac-
214 tion, calibration, flat-field correction, and bad pixels/cosmic rays correction.

215 We extracted the white (1.088 – 1.68 μm) and the spectral (Supplementary Table 1) light
216 curves from the reduced images, taking into account the geometric distortions caused by the tilted
217 detector of the WFC3/IR channel²⁵. The wavelength range of the white light curve corresponds to
218 the edges of the WFC3/G141 throughput (where the throughput drops to 30% of the maximum).
219 In addition, we tested two wavelength grids for the spectral light curves, with a resolving power of
220 20 and 50. We decided to use the latter as it was able to capture the observed water feature more
221 precisely – i.e there were enough data points within the wavelength range of the water feature to
222 produce a statistically significant result.

223 We fitted the light curves using our transit model package *PyLightcurve*, the transit param-
224 eters shown in Supplementary Table 2, and limb-darkening coefficients (Supplementary Table 1)
225 calculated based on the PHOENIX³¹ model, the nonlinear formula³², and the stellar parameters in
226 Supplementary Table 2.

227 More specifically, we fitted the white light curves with a transit model (with the planet-to-star
 228 radius ratio and the mid-transit time being the only free parameters) alongside with a model for the
 229 systematics^{15,25}. It is common for WFC3 exoplanets observations to be affected by two kinds of
 230 time-dependent systematics^{33–36}: the long-term and short-term “ramps”. The first affects each HST
 231 visit and has a linear behaviour, while the second affects each HST orbit and has an exponential
 232 behaviour. The formula we used for the systematics was the following:

$$R_w(t) = n_w^{\text{scan}}(1 - r_a(t - T_0))(1 - r_{b1}e^{-r_{b2}(t-t_o)}) \quad (1)$$

233 where t is time, n_w^{scan} is a normalisation factor, T_0 is the mid-transit time, t_o is the time when each
 234 HST orbit starts, r_a is the slope of a linear systematic trend along each HST visit and (r_{b1}, r_{b2}) are
 235 the coefficients of an exponential systematic trend along each HST orbit. The normalisation factor
 236 we used was changing to n_w^{for} for upwards scanning directions (forward scanning) and to n_w^{rev} for
 237 downwards scanning directions (reverse scanning). The reason for using different normalisation
 238 factors is the slightly different effective exposure time due to the known up-stream/down-stream
 239 effect³⁷. We, also, varied the parameters of the orbit-long exponential ramp for the first orbit in the
 240 analysed time-series (for_{b1}, for_{b2} instead of r_{b1}, r_{b2}), as in many other HST observations the first
 241 orbit was affected in a different way compared to the other orbits⁵. While we used different ramp
 242 parameters from visit to visit, they appear to be consistent, an expected behaviour as the number
 243 of electrons collected per pixel per second is also consistent.

244 At a first stage we fitted the white light curves using the formulas above and the uncertainties
 245 per pixel, as propagated through the data reduction process. However, it is common in HST/WFC3

246 data to have additional scatter that cannot be explained by the ramp model. For this reason, we
247 scaled-up the uncertainties on the individual data points, in order for their median to match the
248 standard deviation of the residuals, and repeated the fitting⁵. From this second step of analysis,
249 we found that the measured mid-transit times were not consistent with the expected ephemeris¹⁹,
250 which we found to be: $P = 32.94007 \pm 0.00003$ days and $T_0 = 2457363.2109 \pm 0.0004$ BJD_{TDB}²⁶.
251 Supplementary Figure 1, shows the difference between the predicted and the observed transit times
252 using the ephemeris in the literature¹⁹ and the one calculated in this work. We used the de-trended
253 and time-aligned – i.e. with TTVs removed – white light curves to also refine the orbital parameters
254 ($a/R_* = 81.3 \pm 1.5$ and $i = 89.56 \pm 0.02$ deg). At a final step, we used all the new parameters
255 (ephemeris, and orbital parameters) to perform a final fit on the white light curves (again having
256 the planet-to-star radius ratio and the mid-transit time being the only free parameters).

257 Supplementary Figure 2 shows the raw white light curves, the detrended white light curves
258 and the fitting residuals as well as a number of diagnostics, while Supplementary Table 3 presents
259 the fitting results. From these, we can see that:

- 260 • the final planet-to-star radius ratio is consistent among the eight different transits, demon-
261 strating the stability of both the instrument and the analysis process,
- 262 • on average, the white light curve residuals show an autocorrelation of 0.17, which is a low
263 number relatively to the currently published observations of transiting exoplanets with HST⁵
264 (up to 0.7), indicating a good fit,
- 265 • uncorrected systematics are still present in the residuals which, on average, show a scatter

266 two times larger than the expected photon noise, and

- 267 • this extra noise component is taken into account by the increased uncertainties, as the re-
268 duced χ^2 is, on average, 1.16.

269 Furthermore, we fitted the spectral light curves with a transit model (with the planet-to-star
270 radius ratio being the only free parameter) alongside with a model for the systematics that included
271 the white light curve (divide-white method¹⁵), and a wavelength-dependent, visit-long slope²⁵:

$$R_\lambda(t) = n_\lambda^{\text{scan}}(1 - \chi_\lambda(t - T_0)) \frac{LC_w}{M_w} \quad (2)$$

272 where χ_λ is the slope of a wavelength-dependent linear systematic trend along each HST visit,
273 LC_w is the white light curve and M_w is the best-fit model for the white light curve. Again, the nor-
274 malisation factor we used was changed to n_λ^{for} for upwards scanning directions (forward scanning)
275 and to n_λ^{rev} for downwards scanning directions (reverse scanning). Also, in the same way as for
276 the white light curves, we performed an initial fit using the pipeline uncertainties and then refitted
277 while scaling these uncertainties up, in order for their median to match the standard deviation of
278 the residuals.

279 Supplementary Figures 3 to 19 show the raw spectral light curves, the detrended spectral light
280 curves and the fitting residuals as well as a number of diagnostics, while Supplementary Table 4
281 presents all the fitting results and average diagnostics per spectral channel. From these, we can see
282 that:

- 283 • the spectral light curves residuals show, on average, standard deviations much closer to the
284 photon noise and lower values of autocorrelation, proving the advantage of using the white
285 light curve as a model compared to the ramp model, and
- 286 • any extra noise component is taken into account by the scaled-up uncertainties, as the re-
287 duced χ^2 is for all channels, on average, close to unity.

288 Finally, the eight spectra of K2-18 b (Supplementary Figure 20) were combined, using a
289 weighted average, to produce the final spectrum (Table 1).

290 **Stellar contamination** K2-18 is a moderately active M2.5V star, with a variability of 1.7% in
291 the B band and 1.38% in the R band⁴⁰. Hence, to make sure that the observed water feature is
292 not the effect of stellar contamination we fitted the observed spectrum with a model that assumes
293 a flat planetary spectrum and contribution only from the star (M2V star as described in Rackham
294 et al. 2018⁴¹). The model that best describes our data has a spot coverage of 26% and a faculae
295 coverage of 73%. We plot this spectrum versus the observed one in Supplementary Figure 21. In
296 addition, we plot the spectrum produced by the spot and faculae combination reported by Rackham
297 et al. 2018⁴¹ and correspond to a 1% I-band variability, for reference. However, as Supplementary
298 Figure 21 shows, the best-fit model cannot describe the observed water feature. From these we
299 conclude that there is no combination of stellar properties that could introduce the observed water
300 feature.

301 **Atmospheric retrieval** We fitted the final planetary spectrum using our Bayesian atmospheric
302 retrieval framework TauREx^{9,27}, which fully maps the correlations between the fitted atmospheric
303 parameters through nested sampling^{38,39}.

304 The atmosphere of K2-18 b was simulated as a plane-parallel atmosphere with pressures
305 ranging from 10^{-4} to 10^6 Pa, sampled uniformly in log-space by 100 atmospheric layers, assuming
306 an isothermal temperature-pressure profile. We initially tested fitting for a number of trace-gases
307 – H_2O ⁴², CO ⁴³, CO_2 ⁴³, CH_4 ⁴⁴ and NH_3 ⁴⁵ – but found only water vapour to play a significant
308 role. Hence, we proceeded only with this molecule. We also included the effect of clouds using a
309 grey/flat-line model, as the quality and wavelength ranges of the currently available observations
310 do not allow us to make any reasonable constraints on the haze properties of the planet. Finally, we
311 included the spectroscopically inactive N_2 as an inactive gas, to account for any unseen absorbers
312 – e.g. methane, which is expected at these temperatures. As free parameters in our models we had:
313 the volume mixing ratio of H_2O (log-uniform prior between 10^{-10} and 1.0), the volume mixing
314 ratio of N_2 (log-uniform prior between 10^{-10} and 1.0), the planetary temperature (uniform prior
315 between 260 and 320 K), the planetary radius (uniform prior between 0.05 and $0.5 R_{\text{Jup}}$, and the
316 cloud top pressure (log-uniform prior between 10^{-3} and 10^7 Pa, where 10^7 Pa represents a cloud
317 free atmosphere). We restricted the temperature prior compared to all the possible temperatures for
318 different values of albedo and emissivity because, since we can detect only water, the temperature
319 of atmospheric part probed must be higher than the freezing point of water (~ 260 K at 1 mbar).

320 We identified three solutions: a) a cloud-free atmosphere containing only H_2O and H_2/He ,

321 b) a cloud-free atmosphere containing H₂O, H₂/He and N₂, and c) a cloudy atmosphere containing
322 only H₂O and H₂/He. The best-fit spectra and the posterior plots are shown in Figure 3. In all
323 cases, a statistically significant atmosphere around K2-18 b was retrieved with an ADI⁵ of 5.0, 4.7
324 and 4.0, respectively. The ADI is the positively defined logarithmic Bayes Factor, where the null
325 hypothesis is a model that contains no active trace gases, Rayleigh scattering or collision induced
326 absorption – i.e. a flat spectrum. An ADI of 5 corresponds to approximately a 3.6σ^{8,9} detection of
327 an atmosphere. The values are too similar to distinguish between the three scenarios.

328 **Data Availability** The data analysed in this work are available through the NASA MAST HST
329 archive (<https://archive.stsci.edu>) programs 13665 and 14682. The molecular line lists used are
330 available from the ExoMol webpage (www.exomol.com). The final and intermediate results (re-
331 duced data, extracted light curves, light curve fitting results and atmospheric fitting results) are
332 available through the UCL-Exoplanets webpage (<https://www.ucl.ac.uk/exoplanets>).

333 **Code Availability** All the software used to produced the presented results are publicly available
334 through the UCL-Exoplanets GitHub page (<https://github.com/ucl-exoplanets>). More specifically,
335 the codes used were:

- 336 • TauREx (https://github.com/ucl-exoplanets/TauREx_public),
- 337 • *Iraclis* (<https://github.com/ucl-exoplanets/Iraclis>), and
- 338 • *PyLightcurve* (<https://github.com/ucl-exoplanets/pylightcurve>).

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