

Modelling the neutrino in terms of Cosserat elasticity

Olga Chervova and Dmitri Vassiliev
(University College London)

17 July 2009

12th Marcel Grossmann Meeting on General Relativity

Subject of talk

Weyl equation (massless Dirac equation):

$$\sigma^\alpha_{\dot{a}b} \nabla_\alpha \xi^b = 0.$$

Here

- σ^α , $\alpha = 0, 1, 2, 3$, are Pauli matrices,
- ∇_α are covariant derivatives with respect to local coordinates x^α , $\alpha = 0, 1, 2, 3$, and
- ξ^b , $b = 1, 2$, is the unknown spinor field.

Will construct an alternative model without spinors, Pauli matrices or covariant derivatives.

Describing a 3D deformable medium

(a) Classical elasticity: displacements only.

(b) Cosserat elasticity: displacements and rotations. See

E. Cosserat and F. Cosserat, *Théorie des Corps Déformables*, 1909. Available from Amazon.

I will assume that there are no displacements, only rotations.

To describe rotations of material points mathematically I attach to each geometric point a *coframe*.

A coframe ϑ in 3D is a triplet ϑ^j , $j = 1, 2, 3$, of orthonormal covector fields. Each covector field ϑ^j has a hidden tensor index: $\vartheta^j = \vartheta^j_\alpha$, $\alpha = 1, 2, 3$.

Same in plain English: a coframe is a field of orthonormal bases.

NB. Coframe lives separately from local coordinates (not aligned with coordinate lines).

The coframe ϑ is an unknown quantity (dynamical variable).

The other dynamical variable is a density ρ .

Choose potential energy from the condition of conformal invariance. Explicit formula:

$$P(x^0) = \int \|T^{\text{ax}}\|^2 \rho dx^1 dx^2 dx^3,$$

$$T^{\text{ax}} = \frac{1}{3}(\vartheta^1 \wedge d\vartheta^1 + \vartheta^2 \wedge d\vartheta^2 + \vartheta^3 \wedge d\vartheta^3).$$

Standard kinetic energy

$$K(x^0) = \int \|\dot{\vartheta}\|^2 \rho dx^1 dx^2 dx^3,$$

$$\dot{\vartheta} = \frac{1}{3}(\vartheta^1 \wedge \partial_0 \vartheta^1 + \vartheta^2 \wedge \partial_0 \vartheta^2 + \vartheta^3 \wedge \partial_0 \vartheta^3).$$

My action (variational functional)

$$S = \int (P(x^0) - K(x^0)) dx^0.$$

Difference with existing models (teleparallelism)

1. I assume metric to be fixed, i.e. I do not vary metric.
2. My Lagrangian has never been considered.

Solving Euler–Lagrange equations

Vary coframe and density to get Euler–Lagrange equations. Too complicated!

Switch to spinors:

coframe ϑ and density $\rho > 0$



nonvanishing spinor field ξ modulo sign

My Lagrangian density $L(\xi)$ is a rational function of ξ , $\bar{\xi}$ and partial derivatives of ξ , $\bar{\xi}$.

Stationary solutions

Look first for stationary solutions

$$\xi(x^0, x^1, x^2, x^3) = e^{-i\varepsilon x^0} \eta(x^1, x^2, x^3), \quad \varepsilon \neq 0.$$

Theorem 1 *In the stationary case my Euler–Lagrange equation is equivalent to a pair of Weyl equations.*

Proof Turns out my Lagrangian factorises as

$$L(\eta) = \frac{L_{\text{Weyl}}^+(\eta) L_{\text{Weyl}}^-(\eta)}{L_{\text{Weyl}}^+(\eta) - L_{\text{Weyl}}^-(\eta)}.$$

Result follows from factorisation. \square

See <http://arxiv.org/abs/0902.1268>.

Question 1. Can I handle the non-stationary case, i.e. when time dependence is arbitrary? Yes, by means of perturbation theory. I look at perturbations of plane wave solutions. Again I get a pair of Weyl equations.

Question 2. Can I make my model relativistically invariant? Yes, by viewing $(1 + 3)$ -dimensional spacetime as a Cosserat continuum. At perturbative level there is no difference between the nonrelativistic and relativistic models.

Question 3. Can I incorporate mass and external electromagnetic field into my model? Yes, by means of a Kaluza–Klein extension, i.e. by adding an extra coordinate x^4 . See <http://arxiv.org/abs/0812.3948>.

Theorem 2 *In the special case with no dependence on x^3 (i.e. for electron in dimension $1 + 2$) the massive version of my model is equivalent to the massive Dirac equation.*

Summary

New mathematical model for fermions.

- Spacetime viewed as Cosserat continuum.
- Lagrangian chosen from condition of conformal invariance.
- Mass and electromagnetic field incorporated via Kaluza–Klein extension.