Numerical MicroLocal Analysis (NMLA)

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"Given frequency domain wave data, the proposed algorithm gives a pointwise estimate of the the number of rays, their slowness vectors and corresponding wavefront curvature. With time domain wave data and assuming the source wavelet is given, the method also estimates the traveltime."



Geometric optics equations / Far field models : Find local asymptotic (large ω) solutions of

$$\frac{\omega^2}{c^2(x)}\hat{u}(x;\omega) + \Delta\hat{u}(x;\omega) = 0$$

 \hat{u} replaced by "ansatz"

$$\hat{u} \simeq \hat{u}^{ray}(x; \omega) = A(x)e^{i\omega\varphi(x)}$$

yields

$$\begin{cases} \|\nabla\varphi(x)\| = \frac{1}{c(x)}\\ 2\nabla\varphi(x) \cdot \nabla A(x) + A(x)\Delta\varphi(x) = 0 \end{cases}$$

For constant medium and far field data, linear (plane wave) phase approximation is a popular choice (beamforming, DOA) :

$$\widehat{u} \simeq \widehat{u}^{ray}(x; \omega) = B(x)e^{i\vec{k}\cdot x}, \qquad \|\vec{k}\| = \frac{\omega}{c}$$

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Rq. : \Leftrightarrow **Plane wave approximation around a point** x_0

$$\varphi(x) \simeq \varphi(x_0) + (x - x_0) \cdot \nabla \varphi(x_0) + \frac{1}{2} (x - x_0)^T H \varphi(x_0) (x - x_0) + \dots$$
yields

$$\widehat{u}(x;\omega) \simeq B(x_0) e^{i\omega(x-x_0)\cdot\nabla\varphi(x_0)}$$

where

$$B(x_0) = A(x_0)e^{i\omega\varphi(x_0)}.$$

The first "general" *N*-rays ansatz we consider is :

$$\widehat{u}(x;\omega) \simeq \sum_{n=1}^{N} B_n(x_0) e^{i\omega(x-x_0) \cdot \nabla \varphi_n(x_0)} \quad x \text{ near } x_0$$

NMLA observable in the frequency domain

The observable data ($\vec{s} \in \{ \|\vec{s}\| = 1 \}$)



$$U_{\alpha}(\vec{s}) = \frac{c(x_0)}{i\omega} \frac{\partial \hat{u}}{\partial r} (x_0 + r\vec{s}; \omega) + \hat{u}(x_0 + r\vec{s}; \omega), \quad r = \frac{\alpha c(x_0)}{\omega}.$$

"should" fit the ansatz form

$$U_{\alpha}(\vec{s}) \simeq \sum_{n=1}^{N} (\vec{s} \cdot \vec{d_n} + 1) B_n e^{i\alpha \vec{s} \cdot \vec{d_n}} , \quad \vec{d_n} = c(x_0) \nabla \varphi_n(x_0)$$

Relaxation towards a linear system

Set

$$\vec{d} = (\cos \theta_{\vec{d}}, \sin \theta_{\vec{d}}).$$

$$U_{\alpha}(\vec{s}) \simeq \int_{0}^{2\pi} (\vec{s} \cdot \vec{d} + 1) \beta(\vec{d}) e^{i\alpha \vec{s} \cdot \vec{d}} d\theta_{\vec{d}}$$
$$\rightarrow U = K_{\alpha} \beta.$$

(Discretization)

$$K_{\alpha,m,n} = (\vec{d}_m \cdot \vec{d}_n + 1) e^{i\alpha \vec{d}_m \cdot \vec{d}_n}, \quad U_m = U_{\alpha}(\vec{d}_m), \quad B_m = \beta(\theta_m).$$

 K_{α} is compact but a regularized inverse is easy to compute (Filter + truncate Fourier modes of β) and its norm is bounded independently of α (Stability).

NMLA filter

$$\beta := \frac{1}{2L(\alpha) + 1} \mathcal{F}^{-1}(\{\widehat{\beta}_{\ell}\}), \quad \widehat{\beta}_{\ell} = H_{\ell} \mathcal{F}(\{U\})_{\ell}$$

with

$$H_{\ell} = D_{\ell}^{-1} \text{ if } |\ell| < L(\alpha)$$
$$= 0 \text{ else.}$$

where

$$D_{\ell}(\boldsymbol{\alpha}) = 2\pi i^{\ell} (J_{\ell}(\boldsymbol{\alpha}) - i \ J_{\ell}'(\boldsymbol{\alpha}))$$

and

$$L(\alpha) = \min\{\alpha, \alpha + \alpha^{1/3} - 2.5\}$$

Chosen such that $\|K'_{\alpha}^{-1}\| < 3 \rightarrow$ (Stability Theorem)

For an exact Plane wave :

If wavefield is a perfect plane wave of direction \vec{d} with amplitude A then

$$U^{plwa}(\vec{s}) = (\vec{d} \cdot \vec{s} + 1) A e^{i\omega\varphi(x_0)} e^{i\alpha \vec{d} \cdot \vec{s}}$$

and

$$\beta^{plwa}(\vec{s}) = A e^{i\omega\varphi(x_0)} S_{\alpha}(\theta_{\vec{d}} - \theta_{\vec{s}})$$

$$S_{\alpha}(\theta) = \frac{\sin((L(\alpha) + \frac{1}{2})\theta)}{(2L(\alpha) + 1)\sin(\frac{\theta}{2})}$$

where $L(\alpha)$ is the (explicitly given) Number of Fourier Modes used to filter β . Fourier modes are also given analytically

$$\widehat{\beta}_{\ell} = A \, e^{i\omega\varphi(x_0)} \, e^{i\ell\theta_{\vec{d}}}$$

Test 2 sources , homogeneous medium NMLA stability. $|\beta(\theta_{\hat{s}})|$ White noise (20%-40%)



Red lines : exact ray angles.

Varying
$$\alpha = \frac{\omega r}{c(x_0)} \simeq L(\alpha)$$
: 10, 20, 50

NMLA 2nd order / near field application ?

- $\alpha = \frac{\omega r}{c(x_0)} \simeq L(\alpha)$ bounds the number of Fourier modes, while we hope to recover dirac masses ...
- Cannot increase α because of the plane wave approximation. $\varphi(x) \simeq \varphi(x_0) + (x - x_0) \cdot \nabla \varphi(x_0) + \frac{1}{2} (x - x_0)^T H \varphi(x_0) (x - x_0) + \dots$ Recall $x - x_0 = \alpha \frac{c(x_0)}{\omega}$, $r = ||x - x_0||$.
- \rightarrow Need to estimate 2nd order terms.

The simplest 2nd order approximation

Assume only one ray in the solution. Constant curvature HF asymptotics



corresponds to the point source solution $H_0^1(\frac{\omega}{c}|x-x_s|)$.

Approximate NMLA data with H_0^1

$$U_{\boldsymbol{\alpha}}(\vec{s}) \simeq \frac{A_0(x_0)}{\left|\frac{i}{4}H_0^{(1)}\left(\frac{\omega}{c}\left|d\,\boldsymbol{d}+r\,\vec{s}\right|\right)\right|} e^{i\omega(\varphi(x_0)-\boldsymbol{d})} \frac{i}{4}H_0^{(1)}\left(\frac{\omega}{c}\left|d\,\vec{d}+r\,\vec{s}\right|\right)$$

 \vec{d} and d yet to be found



 $\theta_{\vec{d}}$ and $\frac{1}{d}$ are the local

ray direction and mean curvature.

Use FMM type asymptotic expansions $\gamma = \frac{\omega d}{c(x_0)}$ is the large parameter. This "new ansatz" yields a curvature correction to the NMLA Fourier modes



Recall : Plane wave approximation Fourier modes were $\hat{a} = i i \psi(a(x_0)) i \ell \theta$

 $\widehat{\beta}_{\ell} \simeq A e^{i\omega\varphi(x_0)} e^{i\ell\theta_{\vec{d}}}$

Test 1 source , homogeneous medium Black : NMLA - Blue : NMLA 2nd order - Red : exact direction.



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Synthetic data numerical simulation (snapshots)



Generated using standard FDTD + ABCs

Source Point localization in Heteogeneous medium

Synthetic data - backward ray tracing using NMLA output (red) versus

Radon (green) and PWD (yellow).

Ray backward propagation and velocity model



Center of red circle is the exact source localization.