

MATHEMATICS 0054 (Analytical Dynamics)  
YEAR 2023–2024, TERM 2

PROBLEM SET #6

This problem set is due at the *beginning* of lecture on Thursday 7 March.

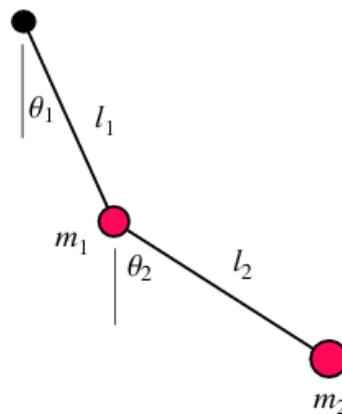
**Topics:** Lagrangian approach to mechanics, continued: Symmetries and conservation laws. Variational principles.

**Reading:**

- Gregory, *Classical Mechanics*, Chapter 12 (handout).
- Handout #11: The Lagrangian approach to mechanics.
- Feynman, *The Feynman Lectures on Physics*, volume 2, Chapter 19 (handout).
- Gregory, *Classical Mechanics*, Chapter 13 (handout).

1. [An old friend: See Problem 3 of Problem Set #3]

A double pendulum consists of rigid massless rods of lengths  $\ell_1$  and  $\ell_2$  and particles of mass  $m_1$  and  $m_2$ , respectively, attached as in the diagram. (All pivots are frictionless, of course.)



- Find the Lagrangian for the system, using as generalized coordinates the angles  $\theta_1$  and  $\theta_2$ .
- Find the exact equations of motion for the system, using the Lagrangian. Do the equations of motion agree with those found by Newtonian methods in Problem Set #3?

2. Let  $F(\mathbf{q}, t)$  be an arbitrary function of the coordinates  $\mathbf{q} = (q_1, \dots, q_n)$  and the time (but *not* of the velocities).

- (a) Show that the Lagrangian  $L'(\mathbf{q}, \dot{\mathbf{q}}, t) \equiv L(\mathbf{q}, \dot{\mathbf{q}}, t) + \frac{d}{dt}F(\mathbf{q}, t)$  leads to the same equations of motion as does the Lagrangian  $L$ . [*Hint:* There are at least two ways of doing this problem. There is a very easy proof using the variational principle; or you can prove it directly by grinding out the derivatives.]

*Remark:* Such a change of Lagrangian is occasionally called a “(Lagrangian) gauge transformation”; though  $L' \neq L$ , the two Lagrangians are physically equivalent, as they lead to the same dynamics.

- (b) Can this be generalized to permit  $F$  to depend on the  $\dot{\mathbf{q}}$  as well?

[There is a partial converse to this theorem: If the Lagrange equations for  $L$  and  $L'$  are *formally identical* — that is, if

$$\Lambda_i(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, t) \stackrel{\text{def}}{=} \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = \frac{\partial^2 L}{\partial \dot{q}_i \partial q_j} \dot{q}_j + \frac{\partial^2 L}{\partial \dot{q}_i \partial \dot{q}_j} \ddot{q}_j + \frac{\partial^2 L}{\partial \dot{q}_i \partial t} - \frac{\partial L}{\partial q_i}$$

is the *same function* of the  $\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}$  and  $t$  as is the analogously defined  $\Lambda'_i(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, t)$  — then  $L'(\mathbf{q}, \dot{\mathbf{q}}, t) \equiv L(\mathbf{q}, \dot{\mathbf{q}}, t) + \frac{d}{dt}F(\mathbf{q}, t)$  for some function  $F(\mathbf{q}, t)$ . For a proof, see Saletan + Cromer, *Theoretical Mechanics*, pp. 40–41.]

3. (a) Show that the Lagrangian function

$$L(\mathbf{r}, \dot{\mathbf{r}}, t) = \frac{1}{2}m\dot{\mathbf{r}}^2 - e\varphi(\mathbf{r}, t) + e\dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

yields the correct equation of motion for a particle with electric charge  $e$  moving in an electromagnetic field, namely

$$m\ddot{\mathbf{r}} = e(\mathbf{E} + \dot{\mathbf{r}} \times \mathbf{B})$$

where

$$\begin{aligned} \mathbf{E} &= -\nabla\varphi - \frac{\partial \mathbf{A}}{\partial t} \\ \mathbf{B} &= \nabla \times \mathbf{A} \end{aligned}$$

are the electric and magnetic fields, respectively. [The vector field  $\mathbf{A}$  is called the *vector potential*, and the scalar field  $\varphi$  is called the *scalar potential*. Both  $\mathbf{A}$  and  $\varphi$  may be functions of  $x, y, z$  and  $t$ .]

- (b) Show that when the potentials of the electromagnetic field are subjected to an “(electromagnetic) gauge transformation”

$$\begin{aligned} \mathbf{A} &\rightarrow \mathbf{A}' \equiv \mathbf{A} + \nabla\psi \\ \varphi &\rightarrow \varphi' \equiv \varphi - \frac{\partial\psi}{\partial t} \end{aligned}$$

where  $\psi(\mathbf{r}, t)$  is an arbitrary function, the electromagnetic field  $\mathbf{E}$  and  $\mathbf{B}$  they describe do not change.

- (c) Determine how the Lagrangian changes if we replace  $\varphi$  by  $\varphi'$  and  $\mathbf{A}$  by  $\mathbf{A}'$ . How is it that the equations of motion are unchanged, despite the fact that  $L' \neq L$ ? (Compare to the preceding problem!)

4. A particle is subject to a constant force  $\mathbf{F}$ .

- (a) Show the Newtonian equations of motion are invariant under spatial translation  $\mathbf{r} \mapsto \mathbf{r}' \equiv \mathbf{r} + \mathbf{e}$ , where  $\mathbf{e}$  is an arbitrary constant vector.
- (b) What does the transformation  $\mathbf{r} \mapsto \mathbf{r}' \equiv \mathbf{r} + \mathbf{e}$  do to the Lagrangian?
- (c) Find the conserved quantity associated with this symmetry by Noether's theorem. What does it express physically?

*Moral:* While translation-invariance of the dynamical law always implies (for a Lagrangian system) the existence of a conserved quantity, that quantity is not always linear momentum.

5. Consider a system of  $N$  point-particles interacting through a potential that depends only on the differences between particle positions, i.e.  $V = V(\mathbf{r}_2 - \mathbf{r}_1, \mathbf{r}_3 - \mathbf{r}_1, \dots, \mathbf{r}_N - \mathbf{r}_1)$ .

- (a) Show that the equations of motion are invariant under a “Galilean boost” with velocity  $\mathbf{u}$ , that is, the transformation  $\mathbf{r}_i \mapsto \mathbf{r}'_i \equiv \mathbf{r}_i + \mathbf{u}t$ . [Cf. the discussion of Galileo's Principle of Relativity in Handout #1.]
- (b) How does the Lagrangian change under a Galilean boost? Show that  $L$  is *not* invariant, but rather undergoes a “(Lagrangian) gauge transformation”

$$L(\mathbf{r}', \dot{\mathbf{r}}') = L(\mathbf{r}, \dot{\mathbf{r}}) + \frac{d}{dt}F(\mathbf{r}, t),$$

and find the function  $F(\mathbf{r}, t)$ .

- (c) Find the constant of motion guaranteed by part (b) and Noether's theorem. What does it express physically?