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### TRANSITION INTEGRALS FOR SITY AND CALL

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ABSTRACT. Atomic wave functions calculated using the Hartree-Fock equation and a polarization potential are used to obtain mean square radii of atomic orbitals for the ions Si<sup>3+</sup> and Ca+, and transition integrals for the SiIV and CaII spectra. A table of radial wave functions for Ca+ is given. Our results are compared with those obtained from other wave functions, and in one case with an experimental value, and generally good agreement is obtained.

- 1. Introduction. In a previous paper (3) by one of us (A.S.D.) a method was described for improving the approximate calculation of energy levels for atomic configurations consisting of a set of closed shells and a single series electron. In the course of that work new radial wave functions were obtained for various states of the ions Si<sup>3+</sup> and Ca<sup>+</sup>. These ions are of considerable astrophysical interest, and we have therefore calculated transition integrals from the wave functions for various spectroscopic transitions (denoted in the customary way as Si IV and Ca II, respectively).
- 2. Results for Si IV. Wave functions were calculated for the ground state  $1s^2 2s^2 2p^6 3s$  of Si<sup>3+</sup> and for the excited states 3p, 3d, 4s, 4p and 4d. Wave functions for these states had previously been calculated including exchange by Hartree, Hartree and Manning (5). The differences between our functions and theirs are quite small. The maximum differences (in the sense Douglas minus Hartree) between the normalized wave functions are: 3s, +0.012; 3p, -0.014; 3d, +0.016; 4s, -0.015; 4p, +0.013; 4d, -0.017. The general effect is a slight contraction of the wave functions, including of course a slight contraction in the positions of the maxima of the wave functions. As a further test of the effect of the introduction of a polarization potential on the wave functions we have computed the mean square radii of the orbitals from our wave functions. The results are given in Table 1, together with values of  $\overline{r^2}$  calculated from

Table 1. Mean square radii in Si<sup>3+</sup>

		$\overline{r^2}$	
Series		s.c.f. with	Coulomb
electron	This paper	exchange	approximation
3s	3.89	3.99	3.78
3p	4.92	5.00	4.68
3d	6.86	7.03	7.03
48	17.0	$17 \cdot 2$	17.0
4p	20.9	21.4	20.8
4d	27.8	29.0	28.8

the Coulomb approximation of Bates and Damgaard (2), footnote, p. 111). The largest change introduced by the present work is about 5 % in  $\overline{r^2}$ . In view of the smallness of these changes we are not publishing the detailed tables of radial wave functions.

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The transition integrals are given in Table 2. The tabulated quantity is  $\sigma^2$ , where

$$\sigma^2 = \frac{1}{4l^2-1} \biggl( \int_0^\infty r P_i P_f dr \biggr)^2, \label{eq:sigma2}$$

l is the greater of the two azimuthal quantum numbers involved in the transition, and  $P_i/r$  and  $P_f/r$  are the normalized radial wave functions. Transition probabilities and oscillator strengths may be obtained from  $\sigma^2$  in the usual way ((1), pp. 51–62). Our results are compared with those given by Hartree, Hartree and Manning (5), and with results we have computed from the Coulomb approximation of Bates and Damgaard (2). The results obtained for the three methods are in very close agreement. This agreement gives some confidence in the reliability of the results, although it does not of course constitute strict proof of their accuracy.

Table 2. Transition integrals for Si IV

		$\sigma^2$	
Transition	This paper	s.c.f. with exchange	Coulomb approximation
3s-3p	1.23	1.26	1.19
3p-3d	0.31	0.32	0.30
3s-4p	0.017*	0.018	0.012
3p-4s	0.33	0.32	0.29
3p-4d	0.0015*	0.00078	0.00073
3d–4 $p$	0.14	0.14	0.14
4s-4p	5.29	5.37	$5 \cdot 20$
4p– $4d$	1.24	1.27	$1 \cdot 25$

<sup>\*</sup> Substantial cancellation occurs.

3. Results for Ca II. Wave functions were calculated for the ground state  $1s^22s^22p^63s^23p^64s$  of Ca<sup>+</sup> and for the excited states 4p and 3d. Wave functions for these states had previously been calculated by Hartree and Hartree (4) with exchange. The differences between our radial functions and those of Hartree and Hartree are rather larger than for  $Si^{3+}$ . Their maximum values are (in the sense Douglas minus Hartree): 4s, -0.036; 4p, +0.029; 3d, -0.044. There is an appreciable contraction of the wave functions and a contraction in the positions of the maxima of the wave functions. The wave functions for Ca+ have become something in the nature of a test case: the 3d wave function in particular is unusually sensitive to the assumptions made in the calculation. The wave functions are given in Table 3. Wave functions, including both exchange and polarization, have also been given by Trefftz and Biermann (10). Their 4s and 4p wave functions are very similar to ours, differing at most by a few units in the third decimal. Their 3d function is appreciably different from ours, being more compact, and differing from ours by up to 0.022. The values of  $\overline{r^2}$  from our wave functions are given in Table 4, together with those from the wave functions of Trefftz and Biermann, and of Hartree and Hartree, and with values from the Coulomb approximation (2), footnote, p. 111). The change in  $\overline{r^2}$  for the 3d electron on introducing polarization is proportionally much greater than for the 4s and 4p electrons.

Table 3. Normalized radial wave functions for Ca+

				•			
r	P(4s)	P(4p)	P(3d)	r	P(4s)	P(4p)	P(3d)
0.01	0.042	0.002	***********	1.4	+0.009	-0.082	0.614
0.02	0.068	0.007	0.0002	1.6	-0.125	+0.005	0.613
0.03	0.081	0.015	0.0006	1.8	-0.250	0.096	0.603
0.04	0.084	0.023	0.0013	$2 \cdot 0$	-0.357	0.182	0.587
0.05	0.079	0.033	0.0024	$2 \cdot 2$	-0.444	0.261	0.565
0.06	0.069	0.043	0.0040	$2 \cdot 4$	-0.510	0.331	0.540
0.07	0.055	0.053	0.0059	$2 \cdot 6$	-0.556	0.390	0.512
0.08	0.038	0.063	0.0083	2.8	-0.585	0.438	0.482
				3.0	-0.599	0.476	0.452
0.10	0.000	0.080	0.0145	$3 \cdot 2$	-0.600	0.504	0.421
0.12	-0.038	0.095	0.022	$3 \cdot 4$	-0.591	0.523	0.392
0.14	-0.072	0.106	0.032	$3 \cdot 6$	-0.575	0.534	0.363
0.16	-0.102	0.113	0.043	3.8	-0.552	0.538	0.334
0.18	-0.124	0.117	0.055	4.0	-0.524	0.535	0.308
0.20	-0.140	0.117	0.069				
0.22	-0.150	0.114	0.083	4.5	-0.446	0.507	0.246
0.24	-0.153	0.108	0.098	$5 \cdot 0$	-0.363	0.459	0.194
0.26	-0.151	0.100	0.114	5.5	-0.289	0.403	0.151
0.28	-0.144	0.090	0.130	$6 \cdot 0$	-0.224	0.344	0.115
0.30	-0.133	0.078	0.147	6.5	-0.170	0.288	0.087
				$7 \cdot 0$	-0.127	0.236	0.066
0.35	-0.092	0.044	0.189				
0.40	-0.038	0.005	0.231	8	-0.069	0.152	0.036
0.45	+0.019	-0.034	0.272	9	-0.035	0.094	0.020
0.50	0.076	-0.072	0.311	10	-0.018	0.056	0.010
0.55	0.128	-0.107	0.348	11	-0.009	0.032	0.005
0.60	0.172	-0.137	0.383	12	-0.004	0.018	0.003
				13	-0.002	0.010	0.002
0.7	0.237	-0.183	0.444	14	-0.001	0.005	0.002
0.8	0.266	-0.208	0.494	15	<del>-</del>	0.003	0.001
0.9	0.265	-0.215	0.533	16		0.001	0.001
1.0	0.239	-0.207	0.563	17	_	0.001	
1.1	0.195	-0.187	0.584				
1.2	0.139	-0.158	0.599				

Table 4. Mean square radii in Ca+

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Series electron	This paper	Trefftz and Biermann	s.c.f. with exchange	Coulomb approximation
48	14.5	14.6	15.8	13.7
<b>4</b> p	21.8	21.7	22.9	$20 \cdot 4$
$\vec{3d}$	6.79	6.33	8.42	6.52

The transition integrals for Ca II spectral transitions are listed in Table 5. Both the dipole length and dipole velocity formulae have been used by Trefftz and Biermann and by ourselves. The values for self-consistent fields with and without exchange are taken from Hartree and Hartree (4) and from Bates and Damgaard (2); Zwann's results (12) are based on the use of an approximate core potential, and the Coulomb approximation tables ((2), together with a normalization correction due to Seaton (9))

had to be judiciously extrapolated to obtain the 4p-3d value. Varsavsky's results (11) are based on Layzer's screening theory (6). The value by Ostrovskii and Penkin (8) is experimental, one of the relatively few absolute intensity measures for lines in the second spectrum of an atom. Finally, we have collected in Table 6 values of the radial integral needed in calculating the probability of the 4s-3d quadrupole transition.

Table 5. Transition integrals for Ca II

		$\sigma^2$	
$\overbrace{4s4p}$		4p	-3d
Length	Velocity	Length	Velocity
4.74	4.45	0.35	0.095
4.76	4.47	0.31	0.11
5.15	4.1	0.46	
6.11	-	1.11	
4.81		0.42	
4.38		0.25	
5.41		0.019	
5.14			
	Length 4.74 4.76 5.15 6.11 4.81 4.38 5.41		

Table 6. Quadrupole integral in Ca II

	$\int_0^\infty r^2 P(4s) \ P(3d) \ dr$	
s.c.f. with exchange	-10.8	
This paper	- 9.1	
Trefftz and Biermann	- 8.6	

Osterbrock (7) used the Hartree and Hartree wave functions for this purpose, and we have quoted his result together with those obtained from our wave functions and from the wave functions of Trefftz and Biermann. As usual it is difficult to assess the reliability of the transition integrals. One is tempted to select a representative value  $\sigma^2 = 4.9$  for the 4s-4p transition. It does seem as if in this case the Coulomb approximation gives a result below the true value. The value of  $\sigma^2$  for the 4p-3d transition is obviously still rather uncertain. An experimental determination of this quantity is greatly desired.

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#### REFERENCES

- (1) Allen, C. W. Astrophysical quantities (London, 1955).
- (2) Bates, D. R. and Damgaard, Agnete. Philos. Trans. Roy. Soc. London. Ser. A, 242 (1949), 101-122.
- (3) Douglas, A. S. Proc. Cambridge Philos. Soc. 52 (1956), 687-692.
- (4) HARTREE, D. R. and HARTREE, W. Proc. Roy. Soc. London. Ser. A, 164 (1938), 167-191.

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(5) HARTREE, W., HARTREE, D. R. and MANNING, M. F. Phys. Rev. 60 (1941), 857-865.

(6) LAYZER, D. Ann. Physics, 8 (1959), 271-296.

(7) OSTERBROCK, D. E. Astrophys. J. 114 (1951), 469-472.

(8) OSTROVSKII, Y. I. and Penkin, N. P. Optika i Spektroskopiya, 10 (1961), 8-14.

(9) SEATON, M. J. Monthly Not. Roy. Astr. Soc. 118 (1958), 504.

(10) TREFFTZ, E. and BIERMANN, L. Z. Astrophys. 30 (1952), 275-281.

(11) VARSAVSKY, C. M. Thesis (Harvard University, 1958).

(12) ZWANN, A. Naturwissenschaften, 17 (1929), 121-122.

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