



*Statistics, Lambda  
and closed  
Universes*



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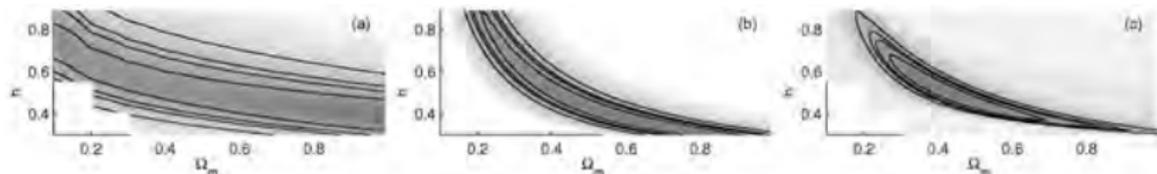
- First met Ofer while giving a graduate course in the CMB
- We then worked with each other as an early example of Cavendish/loA cooperation
- First joint paper was

[arXiv:astro-ph/9802109](https://arxiv.org/abs/astro-ph/9802109) [pdf, ps, other]

## Joint estimation of cosmological parameters from CMB and IRAS data

Matthew Webster, S.L. Bridle, M.P. Hobson, A.N. Lasenby, Ofer Lahav, Graca Rocha

- The idea of this was to break the degeneracies inherent in the CMB alone, or LSS alone, by combining the two

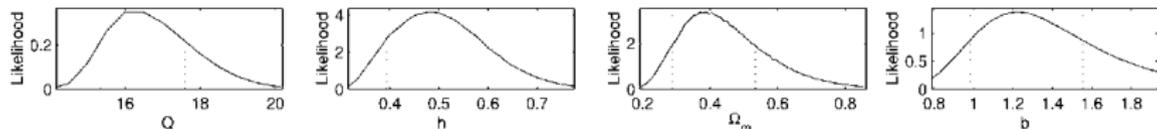


CMB alone

IRAS 1.2 Jy redshift survey

Combined

- The Bayesian statistical framework adopted was pretty much similar to what would be done today - following are 1d marginalised results



## 6. DISCUSSION

The results of this joint optimization are in reasonable agreement with other current estimates. The relatively low value of  $\Omega_m \approx 0.4$  is close to that found by others (White et al. 1993; Bahcall, Fan, & Cen 1997) and is in line with recent supernovae results (Perlmutter et al. 1998). However, given the assumption of a flat universe, it requires a very high cosmological constant ( $\Omega_\Lambda = 0.6$ ).

# Supernovae

- The *Supernovae Cosmology Project* results published July 1997 (Perlmutter *et al.*, *ApJ*, **483**, 565) had found

Comparing light-curve width-corrected magnitudes as a function of redshift of our distant ( $z = 0.35\text{--}0.46$ ) supernovae to those of nearby Type Ia supernovae yields a global measurement of the mass density,  $\Omega_M = 0.88^{+0.69}_{-0.60}$  for a  $\Lambda = 0$  cosmology. For a spatially flat universe (i.e.,  $\Omega_M + \Omega_\Lambda = 1$ ), we find  $\Omega_M = 0.94^{+0.34}_{-0.28}$  or, equivalently, a measurement of the cosmological constant,  $\Omega_\Lambda = 0.06^{+0.28}_{-0.34}$  ( $<0.51$  at the 95% confidence level).

- And constraints on  $\Omega_\Lambda$  from radio-selected gravitational lenses (Falco, Kockhanek & Munoz, *ApJ*, **494**, 47), published February 1998, gave  $\Omega_\Lambda < 0.62$  at  $2\sigma$
- So our joint CMB and LSS result for  $\Lambda$  came in a context where people weren't necessarily expecting a large  $\Omega_\Lambda$ !
- May object that our analysis was limited to flat models
- Most were at that time, but in fact we already had pretty good evidence for high  $\Omega_{\text{tot}}$  from the CMB at the time

## Constraints on cosmological parameters from recent measurements of cosmic microwave background anisotropy

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Accepted 1997 July 31. Received 1997 July 11; in original form 1996 September 19

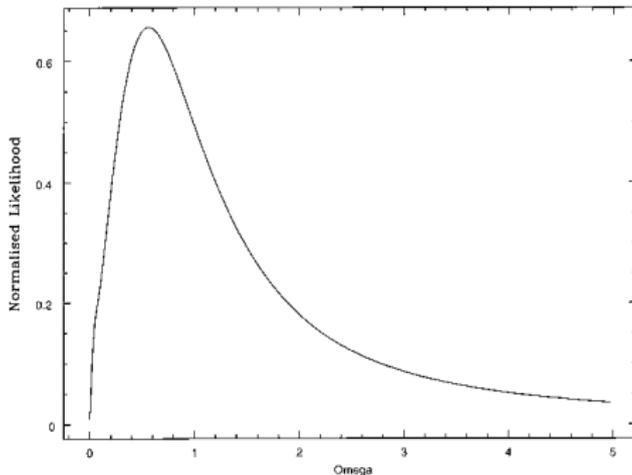
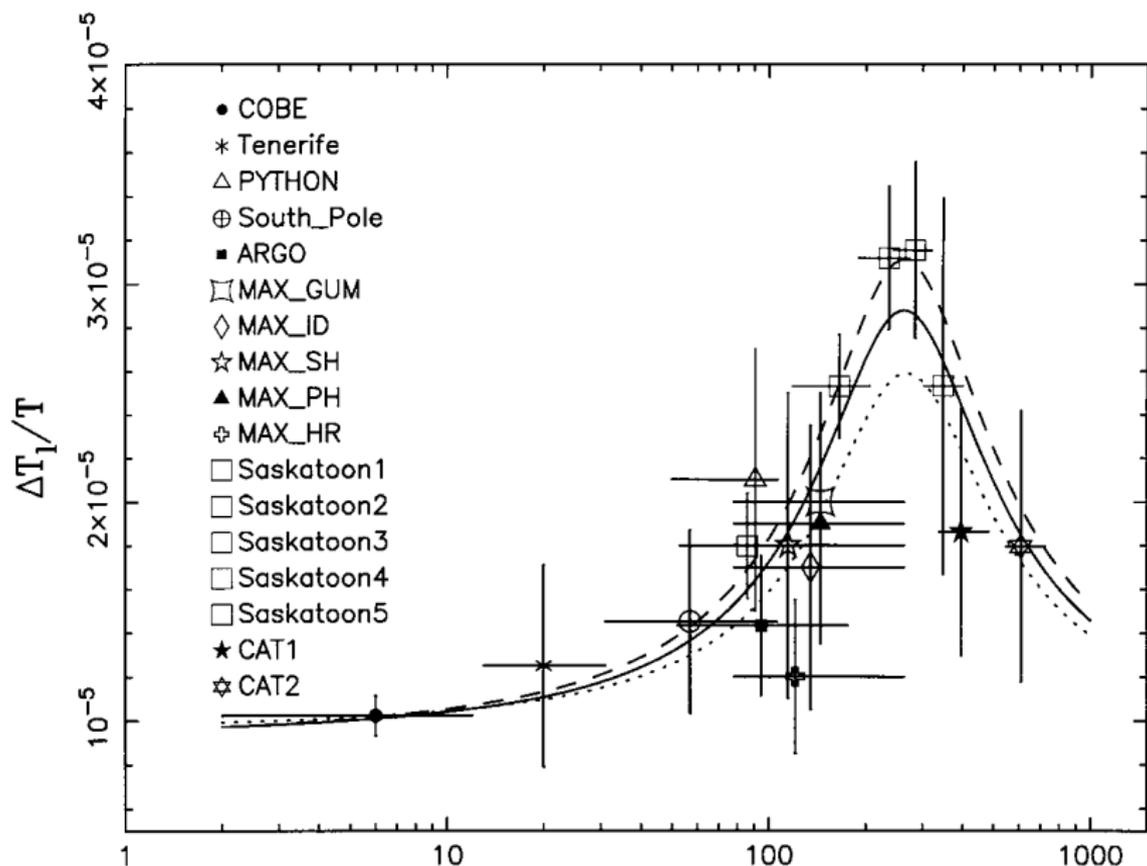


Figure 4. The 1D marginal likelihood curve for  $\Omega$ .

# CMB data ~ 1996



# Hyperparameters

## Bayesian ‘hyper-parameters’ approach to joint estimation: the Hubble constant from CMB measurements

O. Lahav,<sup>1,2★</sup> S. L. Bridle,<sup>3</sup> M. P. Hobson,<sup>3</sup> A. N. Lasenby<sup>3</sup> and L. Sodré, Jr<sup>4</sup>

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Accepted 2000 April 28. Received 2000 April 14; in original form 1999 December 8

We show below that if the prior probabilities for  $\ln(\alpha)$  and  $\ln(\beta)$  are uniform then one should consider the quantity

$$-2 \ln P(\mathbf{w}|D_A, D_B) = N_A \ln(\chi_A^2) + N_B \ln(\chi_B^2) \quad (4)$$

instead of equation (1). Note that in this case our method is equivalent to assuming that we are ignorant of the relative scale of the errors in each experiment.

- This results from generalising the usual equation for combining two data sets  $A$  and  $B$

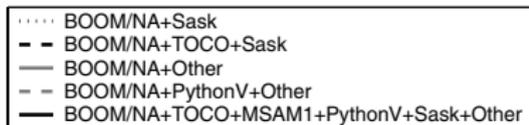
$$\chi_{\text{joint}}^2 = \chi_A^2 + \chi_B^2$$

- We include hyperparameters  $\alpha$  and  $\beta$  in the form

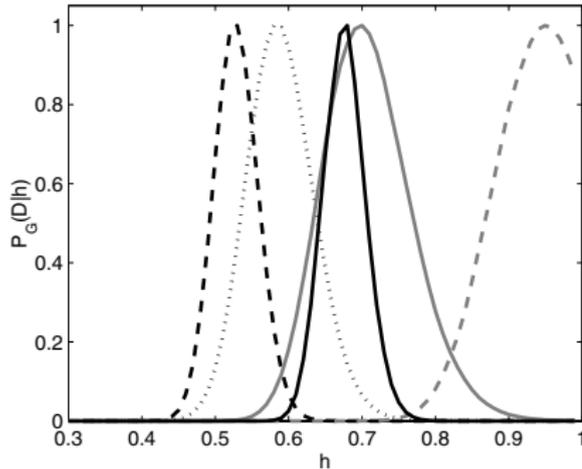
$$\chi_{\text{joint}}^2 = \alpha \chi_A^2 + \beta \chi_B^2$$

which are then marginalised over with a uniform log prior

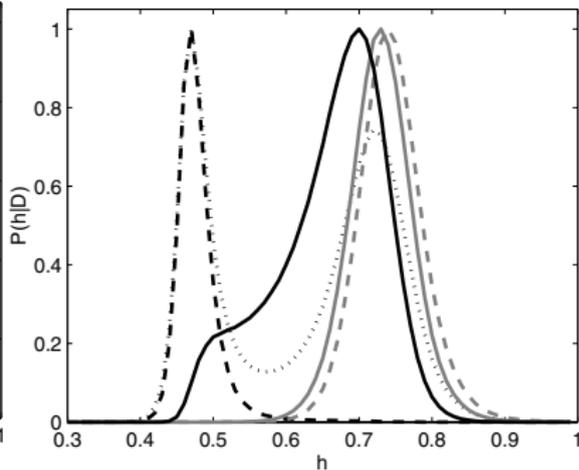
# Hyperparameters — the (year 2000) results for $H_0$



Conventional  $\chi^2$  Analysis



HyperParameter Analysis



## Bayesian sparse reconstruction: a brute-force approach to astronomical imaging and machine learning

Edward Higson <sup>1,2</sup>★ Will Handley,<sup>1,2</sup> Michael Hobson<sup>1</sup> and Anthony Lasenby<sup>1,2</sup>

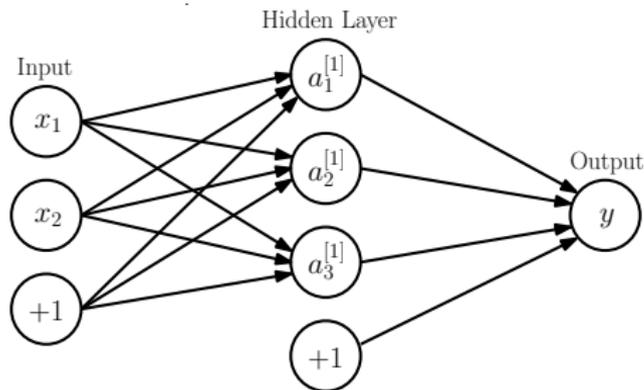
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Accepted 2018 December 3. Received 2018 November 25; in original form 2018 September 10

- Idea here is to explore a fully Bayesian approach to fitting functions and machine learning
- Moreover one in which we can ensure sparsity via choices of prior
- Model a signal as a sum of  $N$  basis functions, of a type  $T$  belonging to discrete classes and with a set of continuous controlling parameters
- Sparsity can be enforced e.g. by a prior on  $N$
- Full exploration of the posterior allows one to carry out the equivalent of ‘dictionary learning’

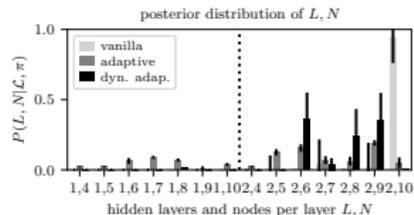
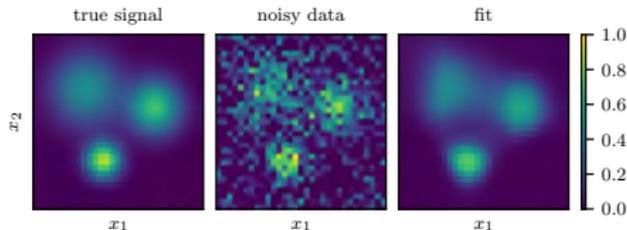
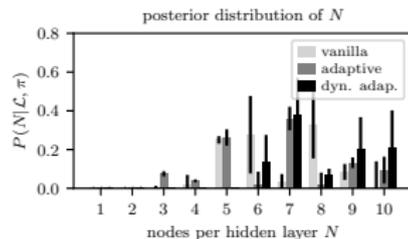
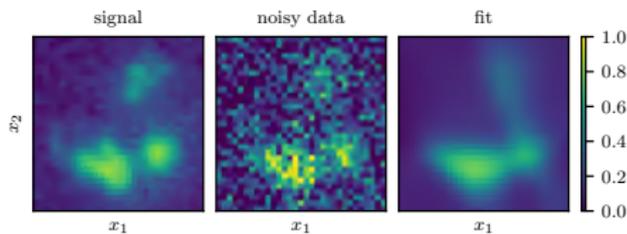
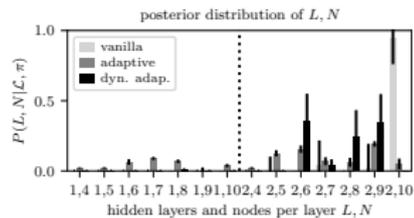
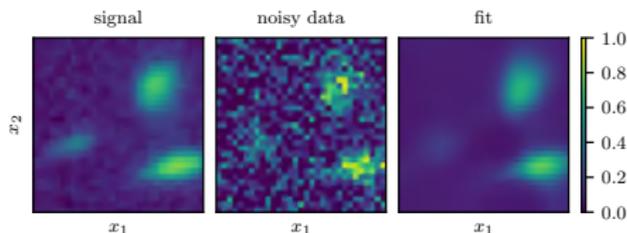
- Implied ‘trans-dimensional’ sampling is dealt with by a product space approach, and ‘Evidence’ not needed as such — we get directly the posterior ratios of different models as subsets of the total space
- This is computationally pretty heavy, so have only applied to toy problems so far
- But quite interesting nevertheless — e.g. one application in paper is to neural network architecture



We can let the **parameters** of this architecture, such as the number of nodes in a hidden layer, be part of the parameters of our problem

Can then ‘marginalise’ over these — quite novel

# 2-D examples: generalised Gaussians & HST images



- Note the computational side of this was made possible by using the new slice-sampling replacement for MultiNest, called *PolyChord* and a *Dynamic Nested Sampling* program called DyPolyChord
- See Handley, Hobson & Lasenby, MNRAS, **450**, L61 (2015), and MNRAS, **453**, 4384 (2015) for PolyChord
- And Higson *et al.* *Astrophysics Source Code Library*, record [ascl:1902.010](#) for the DyPolyChord program (significant speed-ups with both)

# Closed universes

- For some while have been interested in a closed model of the universe satisfying a particular boundary condition
- This is that the total elapse of conformal time, as  $t$  goes from a big bang start to  $t = \infty$ , should be  $\eta = \pi/2$
- Can think of this condition as being 'half' of what a photon would do in pure de Sitter space, but where midpoint of the embedding has been replaced by the big bang

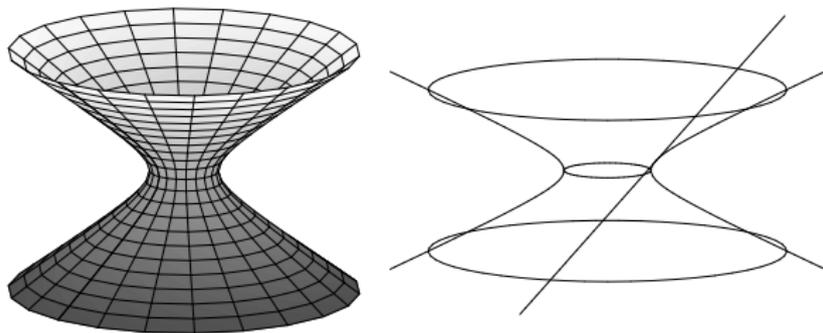
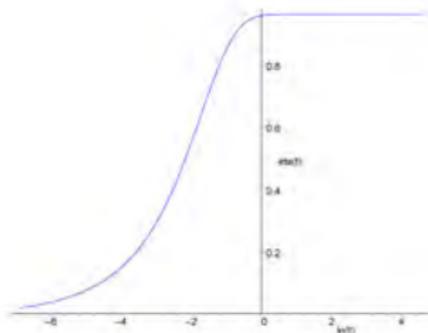


Figure 2: *Two-dimensional de Sitter Space*. The timelike direction is vertical, and spatial sections are closed. The right-hand diagram shows a null geodesic, which is a straight line in the embedding space.

- This very simple requirement (which can be motivated in a variety of alternative ways), when coupled with inflation, has some surprising effects
- This is because inflation ‘eats up’ some conformal time, but then makes it difficult for the universe to achieve more conformal time except by lasting a very long (cosmic) time



- $\eta = \int dt/R(t)$  increases rapidly then saturates at about 0.967
- Therefore need about 0.604 in  $\eta$  to be made up in rest of development of universe

- But the  $R(t)$  by this stage (end of plot) is already  $2.44 \times 10^{22}$ !
- So universe is going to have to last a very long time further to make up the deficit in  $\eta$
- Assume evolution from  $t = t_0 = 14.56$  onwards satisfies a  $R \propto (t - t_0)^{2/3}$  behaviour (crude approx but ok order of magnitude)
- Equation need to solve to find the age  $T$  of the rest of the universe is then:

$$\frac{(100 - 14.56)^{2/3}}{2.44 \times 10^{22}} 3T^{1/3} = 0.604$$

- Therefore need the universe to exist for a time equivalent to  $1.63 \times 10^{61}$  Planck lengths to build up the rest of  $\eta$  (this is about  $2.8 \times 10^{10}$  years).
- This gives a good illustration of how large numbers, such as the age of the universe in Planck times, can come out of inflation plus our boundary condition

- This has implications for the cosmological constant, and in fact can show that

$$\Lambda \approx e^{-6N} l_{\text{pl}}^{-2}$$

- Here  $N$  is the number of e-folds
- $N \approx 46$  (typical for closed universe inflation and nothing to do with our boundary condition) implies  $\Lambda \sim 10^{-122} l_{\text{pl}}^{-2}$ , exactly in the right ball park!
- So all of this was first gone through in the 2005 paper [Lasenby & Doran, Phys.Rev.D 71, 063502](#)
- First full Bayesian comparison with the observations (in which it did quite well!) in [Bridges et al. \(2007\) \(MNRAS, 381, 68\)](#)
- Then further tests in [Vazquez et al., MNRAS 442, 1948 \(2012\)](#) and [JCAP, 08, 001 \(2013\)](#)
- Again did well, but more recently model is in trouble, since based on  $\phi^2$  potential and this now effectively ruled out by current B-mode constraints

- You are probably thinking isn't curvature itself now ruled out?
- Planck 2018 Parameters paper gives results at right:
- But CAMSPEC for same data gives  $\Omega_K = -0.037^{+0.019}_{-0.014}$
- And if add CMB lensing and then BAO get successively  $\Omega_K = -0.0106 \pm 0.0065$  and  $\Omega_K = 0.0007 \pm 0.0019$  though clearly there is some tension underneath this increase in precision

The combination of the *Planck* temperature and polarization power spectra give

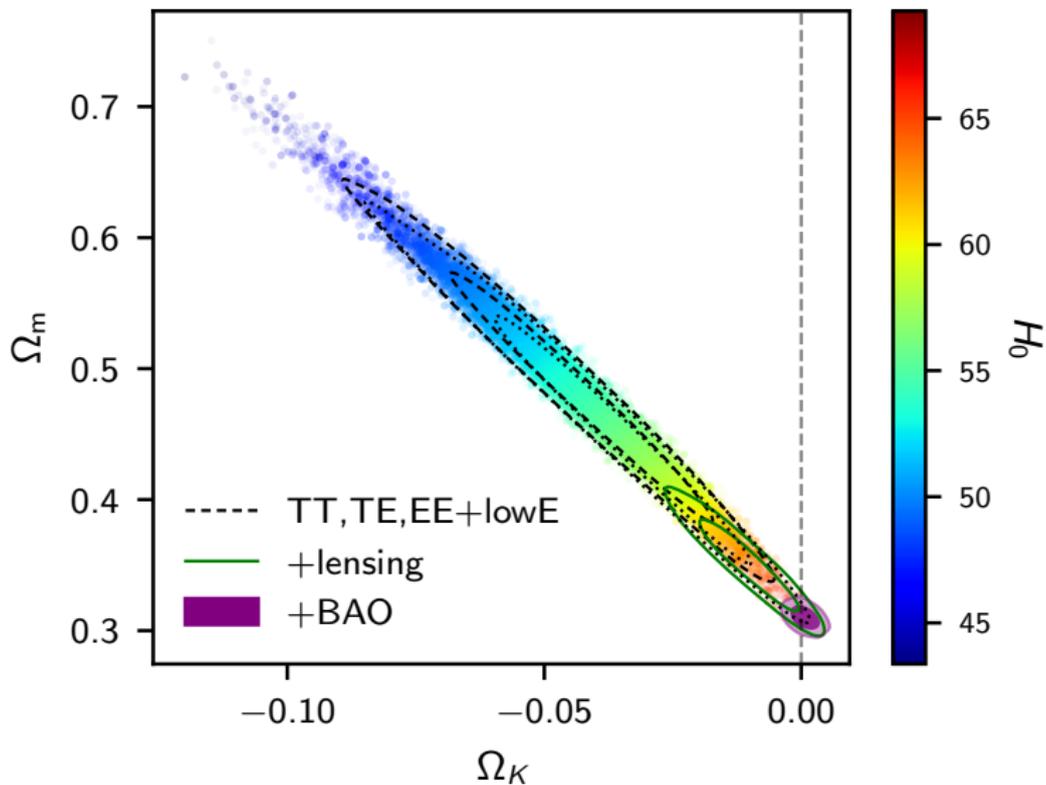
$$\Omega_K = -0.056^{+0.028}_{-0.018} \quad (68\%, \text{Planck TT+lowE}), \quad (46a)$$

$$\Omega_K = -0.044^{+0.018}_{-0.015} \quad (68\%, \text{Planck TT,TE,EE+lowE}), \quad (46b)$$

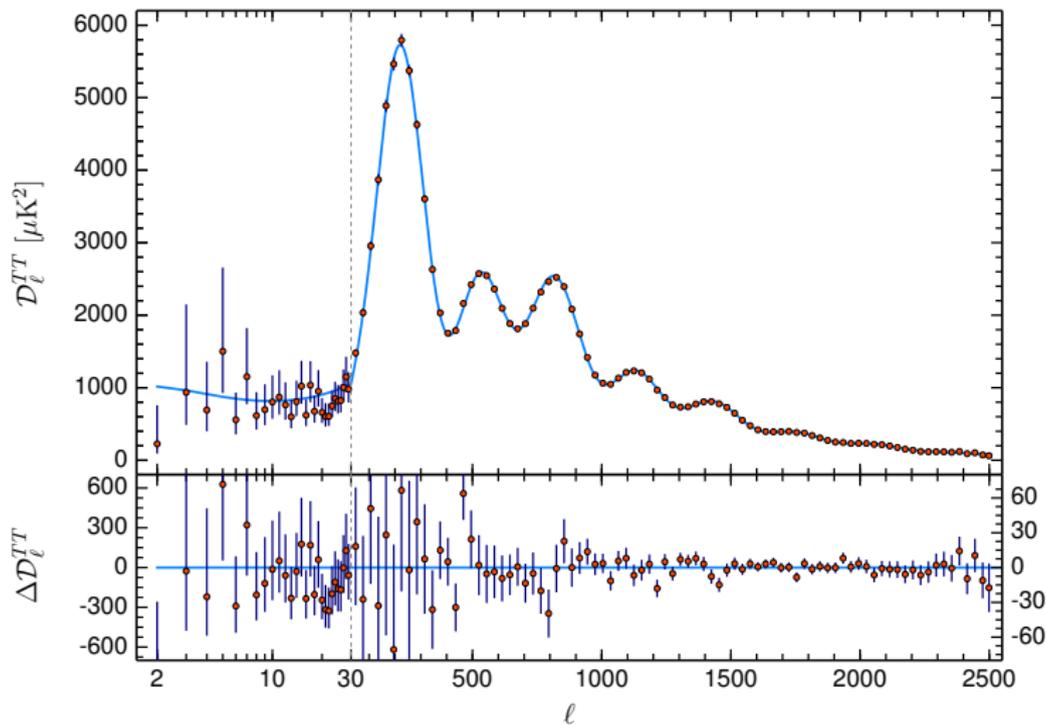
an apparent detection of curvature at well over  $2\sigma$ . The 99% probability region for the TT,TE,EE+lowE result is  $-0.095 < \Omega_K < -0.007$ , with only about 1/10000 samples at  $\Omega_K \geq 0$ .

- In examples below will use what we can call a 'CMB value' of  $\Omega_K = -0.01$

# Planck 2018 curvature constraints

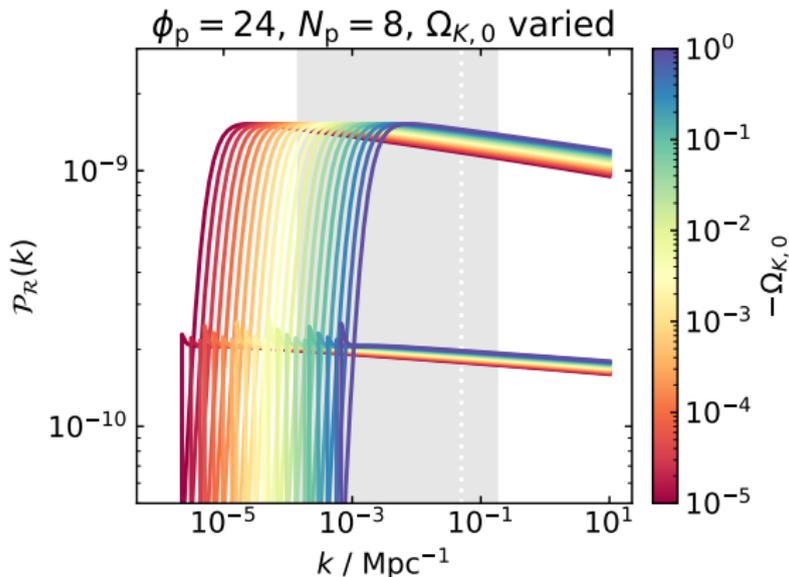


# The 2018 Planck power spectrum



- Generally somewhat low in the low- $\ell$  region and a possible feature around  $\ell = 20 - 30$  — signs of curvature?

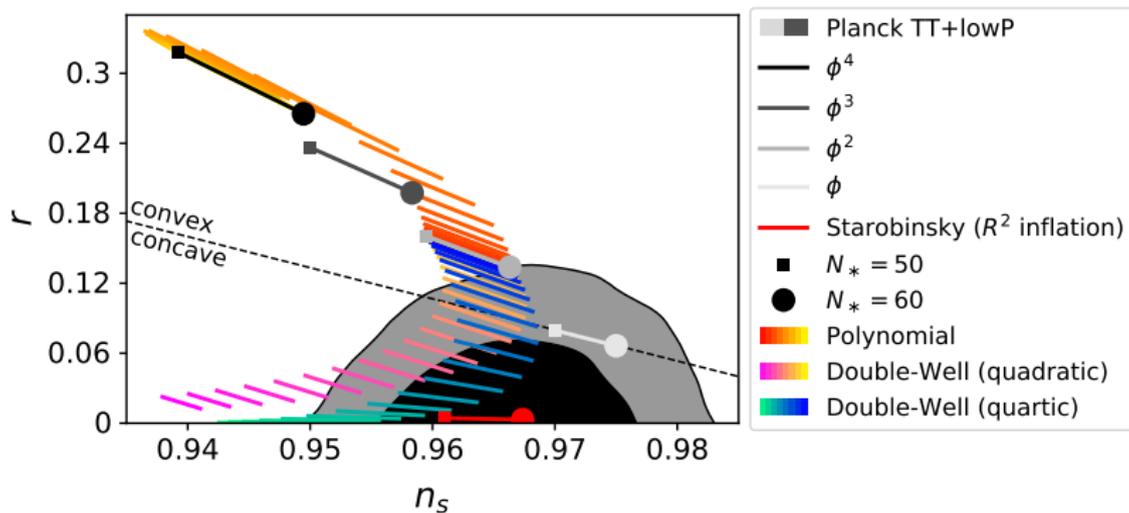
# Cutoff in power spectrum versus $\Omega_k$



(Lukas Hergt *et al.*, in preparation)

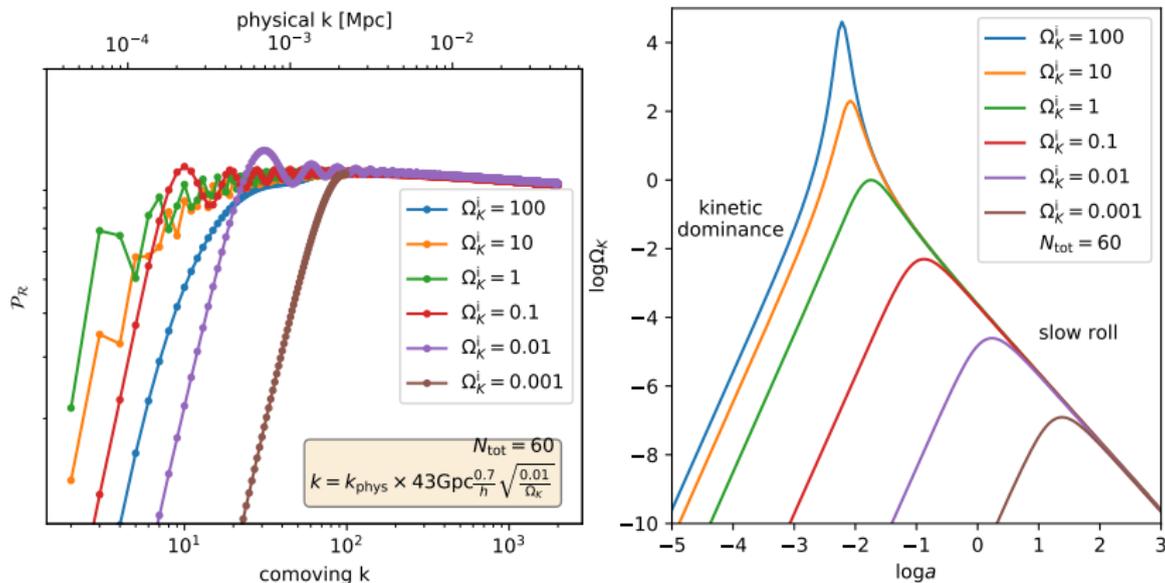
- This computed using slow roll approximation for spectrum and with  $\phi^2$  inflaton potential (same as for previous studies)
- Need to extend to integrating full Mukhanov-Sasaki mode equations and to more general potentials

# $r$ versus $n_s$ for $\phi^2$ and $\phi^4$ combinations



(Lukas Hergt *et al.*, in preparation)

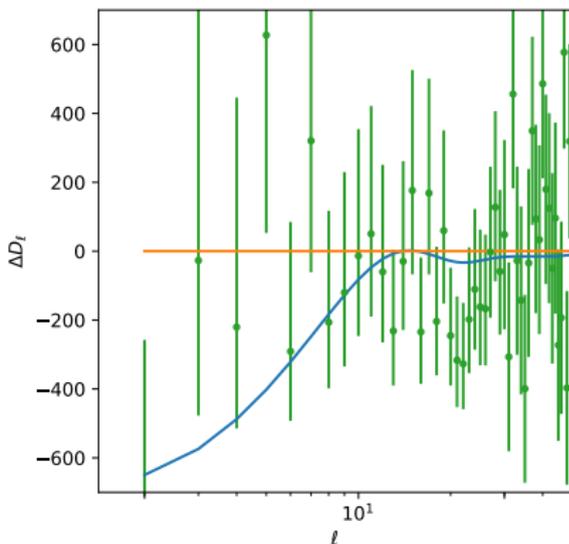
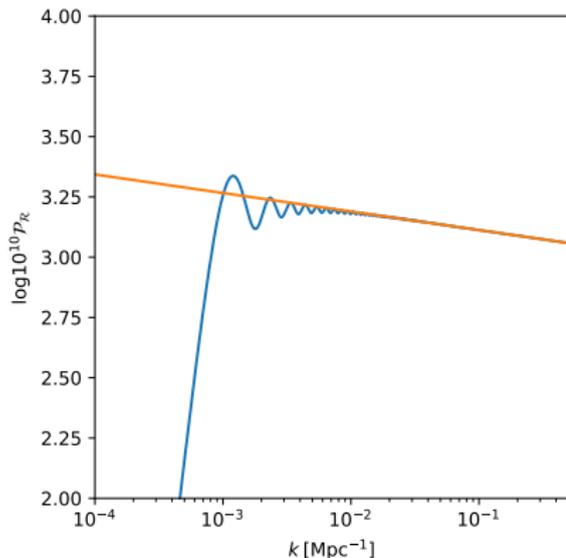
# Primordial power spectrum shape versus curvature



(Will Handley *et al.*, in preparation)

- This uses full Mukhanov-Sasaki mode evolution in closed case (however still for  $\phi^2$  potential)
- Bunch-Davies initial conditions are set at peak curvature

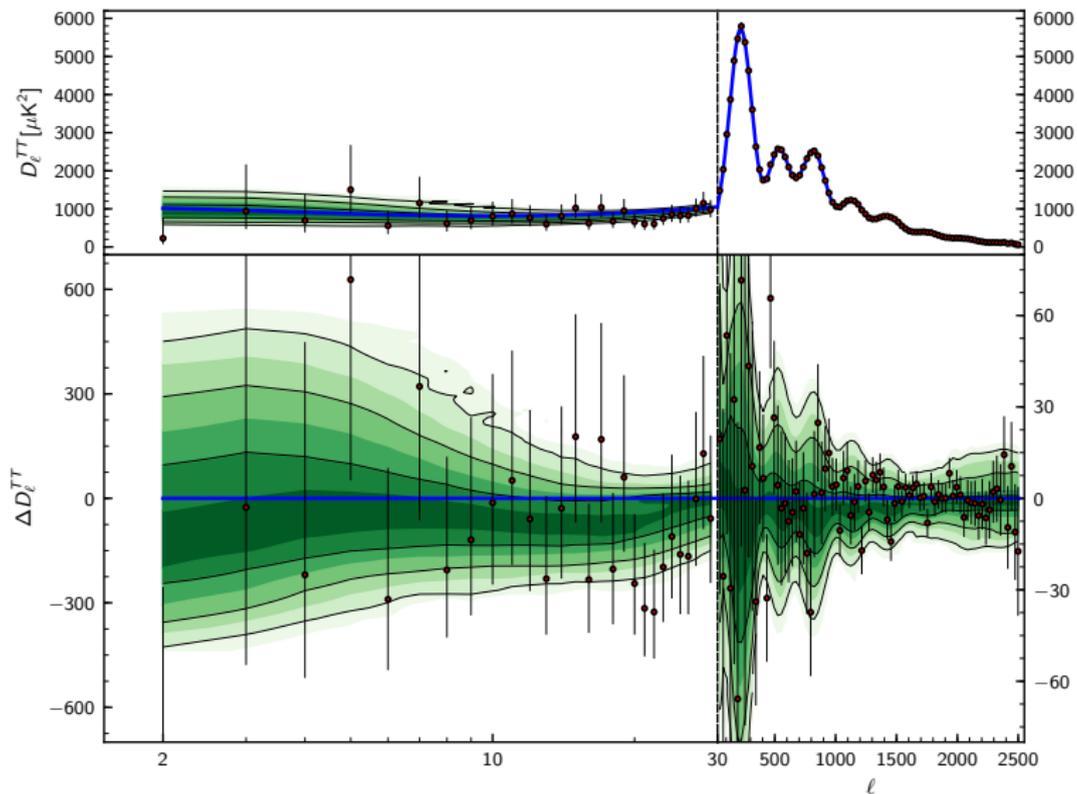
# Effect on CMB power spectrum



(Will Handley *et al.*, in preparation)

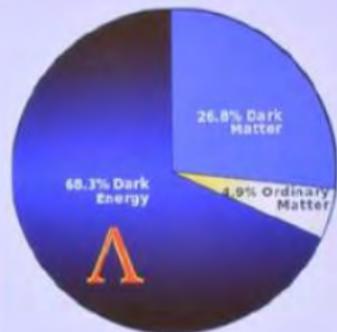
- Worth noting that a free-form reconstruction of the primordial power spectrum from the data, and then transferring *this* through to the  $C_\ell$ s, indicates a lower amplitude of feature favoured in any case

# Effect on CMB power spectrum



(Will Handley *et al.*, in preparation)

## What accelerates the Universe?



*"a simple but strange universe"*