Handout 6 Partial Differential Equations: separation of variables

This is a powerful technique for solving linear PDEs that have no mixed derivatives, i.e. nothing of the form $\partial^2 f / \partial x \partial y$.

There are 6 essential steps:

- 1. Assume the solution is going to be of the form X(x)T(t) or X(x)Y(y), etc. This is called **separable** form.
- 2. Substitute that form back into the PDE.
- 3. Divide by X(x)T(t) or X(x)Y(y).
- 4. Now each term of the equation depends on a different variable so they must both be constants.
- 5. For each possible value of the constant (positive, negative, zero), solve the two resulting ODEs and multiply the solutions together to give one specific solution to the PDE
- 6. Form the general solution of the PDE by adding linear combinations of all the specific solutions.

Example: Heat equation in one dimension

This equation governs the temperature f(x,t) in a thin uniform body of conductivity κ (thin enough that temperature only varies along its length (not across the width):

$$\frac{\partial f}{\partial t} = \kappa \frac{\partial^2 f}{\partial x^2}.$$

Our process goes:

$$f(x,t) = X(x)T(t)$$
 $X(x)T'(t) = \kappa X''(x)T(t)$ $\frac{T'(t)}{T(t)} = \frac{\kappa X''(x)}{X(x)} = \kappa A.$

Zero constant, A = 0

T'(t) = 0 so $T(t) = A_1$ and X''(x) = 0 so $X(x) = B_1x + C_1$, and multiplying these and renaming the constants gives

$$f(x,t) = \alpha x + \beta$$

Negative constant, $A = -\kappa \lambda^2$

 $T'(t) = -\kappa \lambda^2 T(t)$ so $T(t) = A_2 \exp[-\kappa \lambda^2 t]$. $X''(x) = -\kappa \lambda^2 X(x)$ so $X(x) = B_2 \cos(\lambda x) + C_2 \sin(\lambda x)$. Multiplying and renaming again, we get:

$$f(x,t) = \exp\left[-\kappa\lambda^2 t\right] \left(a\cos\left(\lambda x\right) + b\sin\left(\lambda x\right)\right).$$

Positive constant, $A = \kappa \mu^2 \quad T'(t) = \kappa \mu^2 T(t)$ so $T(t) = A_3 \exp[\kappa \mu^2 t]$. $X''(x) = \mu^2 X(x)$ so $X(x) = B_3 \exp[\mu x] + C_3 \exp[-\mu x]$. Putting them together, we have:

$$f(x,t) = \exp\left[\kappa\mu^2 t\right] \left(A \exp\left[\mu x\right] + B \exp\left[-\mu x\right]\right).$$

Note that the temperature here grows exponentially in time: these solutions are not physical!

General solution

$$f(x,t) = \alpha x + \beta + \sum_{n} \exp\left[-\kappa \lambda_{n}^{2} t\right] \left(a_{n} \cos\left(\lambda_{n} x\right) + b_{n} \sin\left(\lambda_{n} x\right)\right) + \sum_{n} \exp\left[\kappa \mu_{n}^{2} t\right] \left(A_{n} \exp\left[\mu_{n} x\right] + B_{n} \exp\left[-\mu_{n} x\right]\right).$$

Real example: heat equation in a finite length bar with cold ends

Now suppose we have a bar of length L which is initially at temperature 1 all over, and which we cool from both ends by holding the ends at temperature 0:

$$f(x,0) = 1$$
 $f(0,t) = 0$ $f(L,t) = 0$

Left hand end Substituting x = 0 into our general solution, we get:

$$f(0,t) = 0 = \beta + \sum_{n} a_n \exp\left[-\kappa \lambda_n^2 t\right] + \sum_{n} \left(A_n + B_n\right) \exp\left[\kappa \mu_n^2 t\right]$$

and forcing this for every possible t gives b = 0, $a_n = 0$ and $A_n + B_n = 0$. Putting these back in makes the full solution become:

$$f(x,t) = \alpha x + \sum_{n} b_n \exp\left[-\kappa \lambda_n^2 t\right] \sin\left(\lambda_n x\right) + \sum_{n} A_n \exp\left[\kappa \mu_n^2 t\right] \left(\exp\left[\mu_n x\right] - \exp\left[-\mu_n x\right]\right)$$

Right hand end Next we look at the condition f = 0 at x = L. This gives:

$$0 = \alpha L + \sum_{n} b_n \exp\left[-\kappa \lambda_n^2 t\right] \sin\left(\lambda_n L\right) + \sum_{n} A_n \exp\left[\kappa \mu_n^2 t\right] (\exp\left[\mu_n L\right] - \exp\left[-\mu_n L\right]).$$

Again, this has to be true for all values of t and the t-dependence of each term is different: so we end up with $\alpha = 0$, $A_n = 0$ and $\lambda_n L = n\pi/L$.

Now almost all the terms have disappeared and the full solution becomes:

$$f(x,t) = \sum_{n} b_n \exp\left[-\kappa n^2 \pi^2 t/L^2\right] \sin\left(n\pi x/L\right).$$

Initial condition Now we put in t = 0 and get:

$$1 = \sum_{n} b_n \sin\left(\frac{n\pi x}{L}\right)$$

which is a Fourier sine series with period 2L that we've seen before:

$$b_n = \frac{2}{L} \int_0^L \sin\left(\frac{n\pi x}{L}\right) dx = \frac{2}{L} \left[-\frac{L}{n\pi} \cos\left(\frac{n\pi x}{L}\right)\right]_0^L = \frac{2}{n\pi} \left(1 - \cos\left(n\pi\right)\right) = \begin{cases} 0 & n \text{ even} \\ \frac{4}{n\pi} & n \text{ odd} \end{cases}$$

We can put this back into our general solution:

$$f(x,t) = \sum_{n \text{ odd}} \frac{4}{n\pi} \exp\left[-\frac{\kappa n^2 \pi^2 t}{L^2}\right] \sin\left(\frac{n\pi x}{L}\right).$$

Let's plot this function for a few values of t:

